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Gas-phase metallicity profiles of the Bluedisk galaxies: Is metallicity in a local star formation regulated equilibrium?

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ABSTRACT

As part of the Bluedisk survey, we analyse the radial gas-phase metallicity profiles of 50 late-type galaxies. We compare the metallicity profiles of a sample of H I-rich galaxies against a control sample of H I-‘normal’ galaxies. We find the metallicity gradient of a galaxy to be strongly correlated with its H I mass fraction ($M(\text{H I})/M_*$). We note that some galaxies exhibit a steeper metallicity profile in the outer disc than in the inner disc. These galaxies are found in both the H I-rich and control samples. This contradicts a previous indication that these outer drops are exclusive to H I-rich galaxies. These effects are not driven by bars, although we do find some indication that barred galaxies have flatter metallicity profiles. By applying a simple analytical model, we are able to account for the variety of metallicity profiles that the two samples present. The success of this model implies that the metallicity in these isolated galaxies may be in a local equilibrium, regulated by star formation. This insight could provide an explanation of the observed local mass–metallicity relation.

Key words: galaxies: abundances – galaxies: evolution – galaxies: ISM.

1 INTRODUCTION

Galaxy formation has been much studied over the past decades, but despite significant successes in this endeavour, it remains unclear exactly how disc galaxies evolve at late times. A particular stumbling block has been determining exactly through which processes a galaxy acquires its cold gas. One of the current leading scenarios is the accretion of gas into the galaxy halo. This halo gas subsequently cools to form a gas disc, from which stars will form (White & Rees 1978; Fall & Efstathiou 1980; Mo, Mao & White 1998). It is predicted that the angular momentum of the accreting gas will grow over time. The angular momentum of the gas disc will therefore also increase, and thus gas cooling from this will settle at increasing radii. This paradigm is commonly referred to as ‘inside-out’ growth (Pichon et al. 2011; Stewart et al. 2013).

The study of the formation and evolution of disc galaxies is complicated by the complex nature of star formation and the cycle of gas within the interstellar medium (ISM) of galaxies. Nevertheless, with three fundamental observables, namely stellar mass, gas-phase

metallicity¹ and the star formation rate (SFR), we can begin to unravel the life of galaxies. The gas-phase metallicity, herein simply referred to as metallicity, is of particular interest since it is not simply a result of star formation integrated through time, but it is also strongly affected by gas flows into and out from the galaxy. With the advent of large spectroscopic fibre surveys, such as the Sloan Digital Sky Survey (SDSS; York et al. 2000) and Galaxy And Mass Assembly project (Driver et al. 2011), a host of studies have explored the relationships between these three aforementioned parameters. Of particular interest is the mass–metallicity relation (Tremonti et al. 2004; Foster et al. 2012), which shows the most massive galaxies to be also the most metal rich. This correlation is commonly attributed to either a downsizing scenario, whereby the most massive galaxies are more efficiently forming stars, or alternatively that galactic-scale winds are more effective at expelling metals from lower mass galaxies. For a more in depth discussion of these and other mechanisms, we refer the reader to Ellison et al. (2008).

¹ By gas-phase metallicity, we refer to the oxygen abundance ($12 + \log_{10}(\text{O}/\text{H})$) of the ISM.

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More recently, studies have reported a secondary correlation of the mass–metallicity relation with the SFR, forming the so-called fundamental metallicity relation (FMR; Mannucci et al. 2010). The FMR presents an anticorrelation of metallicity with the SFR, which has been attributed to either inflows that dilute metallicity and/or outflows that remove metals. We hasten to point out, however, that the FMR is not without contest, with some debate over its origin or existence (Sánchez et al. 2013; Wuyts et al. 2014).

While we may have copious measurements for the central metallicities of galaxies, comparatively less well studied are radial trends of metallicity in galaxies. Early on the picture emerged that at late times ($z \lesssim 0.1$) disc galaxies all show negative (declining radially outwards) metallicity gradients, and when expressed in terms of optical scale radii they showed remarkably similar gradients (Vila-Costas & Edmunds 1992; Zaritsky, Kennicutt & Huchra 1994). Recently Ho et al. (2015) have shown that this common metallicity gradient can be explained by the coevolution of gas, metals and stars.

The common metallicity gradient, however, only applies to isolated galaxies. In a study of interacting systems, Rich et al. (2012) showed a clear tendency towards flatter metallicity gradients, which for the early stages of interaction is consistent with simulations (Rupke, Kewley & Barnes 2010; Torrey et al. 2012).

Despite this work, it is only recently that metallicity gradients have been systematically determined for large samples of galaxies (Moran et al. 2012; Sánchez et al. 2013). With the good statistics these studies provided, these authors have shown the existence of a correlation between stellar mass density and metallicity. This correlation is commonly referred to the local mass–metallicity relation and, as with its global counterpart, its origin is unclear.

In the work of Moran et al. (2012), whose galaxies formed part of the GALEX Arecibo SDSS Survey (Catinella et al. 2010), they attempt to connect metallicity to the atomic gas (H I) content of galaxies. They show tentative hints that the most H I-rich galaxies exhibit sudden drops in metallicity in their outer discs. It is these hints that provided impetus for the work we present here.

In this paper, we present resolved metallicity profiles for all 50 low-redshift ($z \sim 0.025$) galaxies that form the Bluedisk survey (Wang et al. 2013, herein Paper I). The resolved H I maps are the cornerstone of Bluedisk project, providing both the structure and kinematics of the atomic gas disc. The goal of the Bluedisk project is to study in detail two classes of galaxies: an ‘H I-rich’ sample, consisting of those galaxies with stellar masses above $10^{10} M_{\odot}$ and with excess atomic gas, and for comparison a ‘control’ sample consisting of galaxies of similar stellar mass, whose H I content is normal or mildly poor.

We structure this paper as follows: in Section 2, we outline the existing data of the Bluedisk galaxies. In Section 3, we describe our observations and discuss our data-reduction process. In Section 4, we detail our spectral fitting procedures, discuss the global properties of our galaxy population. We also explain our method for determination of metallicity, among other quantities. In Section 5, we present our results, focusing on the resolved metallicity of the Bluedisk galaxies. We use Section 6 to develop and apply a simple model to explain the radial metallicity profiles in terms of their gas and stellar mass contents. Finally, we provide our concluding remarks in Section 7. Throughout this paper, we assume a cosmology with $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_m = 0.3$ and $\Omega_{\Lambda} = 0.7$.

2 DATA

Measurements of the H I content of the Bluedisk galaxies have been obtained using the Westerbork Synthesis Radio Telescope (WSRT) with observations for 49 out of the sample of 50 galaxies, including one non-detection. SDSS images have been analysed to provide optical properties of the galaxies. A full description of the analysis is available in Paper I. We will make use of this data in the context of our new optical spectroscopic data.

2.1 Bluedisk galaxy classification

A key aspect of the Bluedisk strategy is the classification of galaxies into two well-matched H I-rich and control samples. As uncovered by Catinella et al. (2010) there exists a scaling relation between the H I mass fraction ($f_{\text{H I}} = M_{\text{H I}}/M_{*}$), the stellar mass surface density and the observed NUV – r colour. Using the difference between the observed and expected H I mass fractions, we can bisect the Bluedisk population into H I-rich and control samples. We require the samples to contain only isolated galaxies, thus an additional category of non-isolated galaxies has been formed, namely the ‘excluded’ sample. All three samples are described at length in Paper I. We make, however, one minor modification to the classifications listed therein, such that in this paper we consider that BD 39, formerly part of the control sample, to now be a member of the excluded sample. Our motivation for this being that we have identified potentially a small neighbouring galaxy at the same redshift. As a result, this leaves 23 H I-rich galaxies and 18 control galaxies, which we will focus our attention on.

2.2 Bulge–Disc–Bar decomposition

We determine the inclination of the stellar disc using SDSS r -band images. With these images we also perform a bulge–disc–bar decomposition, enabling use to measure a bulge-to-total (B/T) ratio for each galaxy.

Our procedure for bulge–disc–bar decomposition mostly follows that of Weinzirl et al. (2009). The steps are as follows.

(i) We use the SExtractor software (Bertin & Arnouts 1996) to measure the position, ellipticity (e) and position angle (PA) for each galaxy. We also make mask images that flag all neighbouring sources using the segmentation map produced by SExtractor. These masks are used in all the following isophote and model fitting steps.

(ii) Using the photometric measurements as an initial guess, we perform isophote ellipse fitting on the images, and obtain surface density, PA and e profiles for the galaxies. With the surface density profile, we accurately measure the background surface density of the images. The shape (size, PA and e) of the outermost isophote is viewed as the shape of the galactic disc. See Wang et al. (2012) for further details of this step.

(iii) We use the GALFIT package (Peng et al. 2002) to fit models of the bulge, disc and bar to the galaxies. We use exponential models to represent the discs, and use Sérsic models to represent bulges and bars. The Sérsic index is allowed to vary within 1.2–6 for bulge models and vary within 0.3–0.9 for bar models. These are typical values reported in Gadotti & Kauffmann (2009). We first fit a single bulge model, followed by a bulge+disc model and finally bulge+disc+bar model. The model obtained from each fitting step is used as initial guess for the following fitting step. When fitting discs and bars, the size, PA and e measurements from step 1 are used as initial guess. During the fitting, the PA of bars and discs

and are allowed to vary in a very small range ($\pm 20^\circ$). The e of the disc is fixed, and the e of bar is required to be smaller than the e of disc. So in the end, we have three types of models (bulge, bulge+disc and bulge+disc+bar) for each galaxy. We choose the model with the minimum value of reduced χ^2 calculated from the residual map as the best model.

2.3 Stellar mass densities

To derive spatial resolved maps of the stellar mass density, Σ_* , we use SDSS *ugriz* photometry.

We fit the five-band SDSS photometry using the composite stellar population synthesis (SPS) models of BC03, applying the procedure described in Kauffmann et al. (2003a) and Gallazzi et al. (2005). The SPS models combine an underlying exponentially declining star formation history with random bursts of star formation superposed on this. The modelling includes a dust component. The flux from young stars (< 10 Myr) is attenuated following a dust attenuation curve of the form $\tau(\lambda) \propto \lambda^{-1.3}$. Whereas, the flux from long-lived stars (> 10 Myr) is attenuated by a $\tau(\lambda) \propto \lambda^{-0.7}$ power law. The library is described in more detail in Gallazzi et al. (2005). From the results of this SPS fitting, we obtain a posterior distribution on Σ_* .

To perform this SPS modelling, we require a good signal-to-noise (S/N) across the images. We achieve this by using the weighted Voronoi tessellation method of Diehl & Statler (2006), a generalization of the algorithm by Cappellari & Copin (2003). We define our measure of S/N from that of the $u - z$ colour maps (the $u - z$ combination typically offering the poorest S/N). Adopting a threshold S/N = 5, we therefore ensure a good S/N in all colour maps. As the SDSS images contain foreground objects, we mask these objects manually, along with any other spurious features.

3 OBSERVATIONS

Optical long-slit spectroscopic observations of all 50 Bluedisk galaxies were performed in 2013 January and May using the Intermediate dispersion Spectrograph and Imaging System (ISIS) on the 4.2 m William Herschel Telescope (WHT), in a variety of seeing conditions (0.7–1.7 arcsec FWHM). The ISIS spectrograph was operated in a dual arm mode using the standard 5300 Å dichroic, with the GG495 blocking filter in the red arm. Employing the R600B and R600R gratings in the blue and red arms, respectively, a discontinuous spectral coverage of 3700–5300, 5750–7200 Å was provided, with a spectral resolution of ~ 1.7 Å FWHM (full width at half-maximum) constant across all wavelengths. Each target was observed with a minimum of 3×1200 s exposures. A slit width of 3 arcsec was used for all observations presented here, optimizing emission line S/N at the expense of spectral resolution.

Each spectroscopic slit was positioned to coincide with the centre of the galaxy, as defined by the SDSS photometric catalogue. The orientation of the spectroscopic slit was aligned to the kinematic major axis determined from the WSRT velocity moment maps. Where this angle was close to that of a clear natural axis of the galaxy, the slit was more precisely aligned with this optical feature. A final design requirement was to ensure a region observed by a 3 arcsec fibre in the SDSS spectroscopic catalogue was included in the slit, which for all but one target (BD 31), was at the galaxy centre. Overall, the general result was such that the slits were aligned with the optical major axis.

Standard bias frames were obtained for each night, in addition to lamp flat and twilight sky flat exposures. After each target pointing,

additional spectroscopic calibration images were obtained, which comprised Cu-Ar+Cu-Ne arc-lamp exposures, as well as an observation of a standard star. The standard stars were selected from the ING spectrophotometric catalogue² and were observed at a similar airmass to the targeted galaxy. Finally, for characterization of the charge-coupled devices (CCDs), a set of dark frames were also gathered.

3.1 Reduction

Standard PYRAF tasks were used to calculate the CCD bias offsets, pixel gain variations and telescope vignetting effects. With the imcombine task cosmic ray rejection was performed using a mean image combine and a $+3\sigma$ rejection. Bad-pixel masks were constructed manually, based upon the dark and lamp-flat frames. Wavelength calibration was performed using a custom routine, which fitted the arc-lamp spectra at multiple points along the spatial axis. With the use a 2D-spline interpolation, the corresponding wavelength of every pixel was identified. Subsequent sky subtraction was carried out using blank regions in slit.

The dispersion axis of the spectrograph was not perfectly aligned with the CCD axes, this rotation, while small, induced a small ($\lesssim 4$ arcsec) shift in spatial position in the spatial CCD coordinates between the wavelength extrema. To remove this effect, the centre of the target was traced by binning the spectrum in numerous wavelength bins, and fitting a symmetric profile to the spatial intensity distribution in each bin. Sérsic and Gaussian profiles were assumed for galaxy and standard star targets, respectively. A linear fit to this produced a mapping of wavelength to object centre, from which a rectified 2D spectrum was created.

Flux calibration was performed in a two-step procedure. First, an absolute flux calibration by comparing the response of the standard star against the reference spectrum. This was refined by extracting a spectrum from a 3 arcsec² effective aperture matched to the equivalent position of the SDSS fibre spectrum (3 arcsec diameter). We additionally applied a telluric correction by measuring the transmission of the standard star, adopting a linear pseudo-continuum across the affected spectral regions.

On inspection of the data, spectral information was significantly convolved with the spatial profile of the slit, this is an expected consequence of using a spectroscopic slit wider than the seeing disc. For clumpy emission line regions, asymmetric line profiles will be produced and as a result, erroneous velocities will be inferred. Worthy of note is the optical design of the ISIS instrument, which has the dispersion axes in opposite directions for the red and blue arms. This results in ‘contrary offsets’ in the two arms, as shown in Fig. 1. We must properly account for these in our analysis.

4 ANALYSIS

A standard approach to emission line modelling is to assume the lines can each be approximated by a single Gaussian function all with the same velocity offset and dispersion. We perform this spectral fitting using the SDSS PLATEFIT spectral fitting routine (Tremonti et al. 2004; Brinchmann et al. 2004), which first fits a continuum to the spectrum with the emission line features masked, before fitting a sum of Gaussian functions to the residual spectrum. The velocity offsets of the continuum and emission line components are not tied together. The velocity of the emission line component may

² <http://catserver.ing.iac.es/landscape/tm065-100/workflux.php>

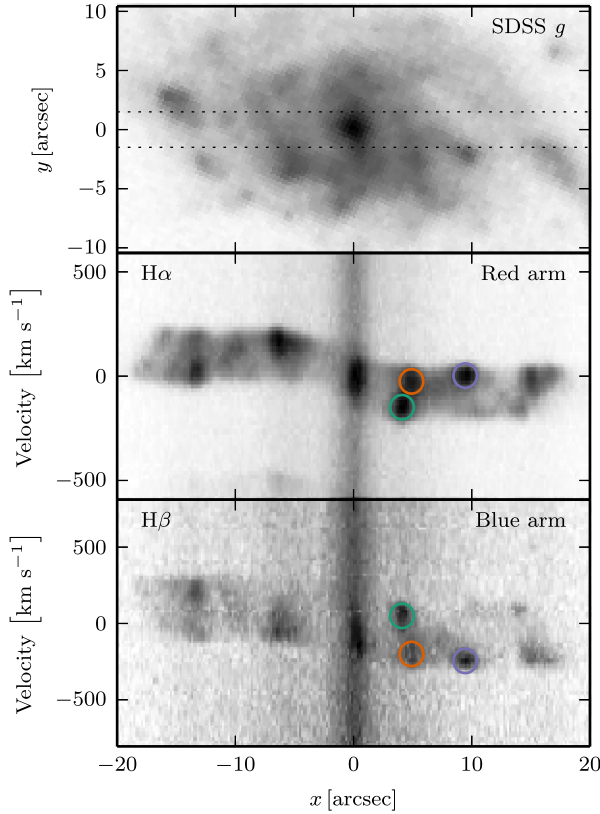


Figure 1. Illustration the ‘contrary offset’ effect of the ISIS instrument, due to the opposite alignment of the dispersion axes of the red and blue spectral arms. Top: SDSS *g*-band image of BD 20 with 3 arcsec wide slit indicated by the dotted horizontal lines. Middle and Bottom: corresponding 2D spectra centred about the $H\alpha$ and $H\beta$ emission line features of the red and blue arms. While the overall velocity curve is preserved, the vertical position in the slit, i.e. in the y -direction, of emitting region is convolved with the velocity information. The coloured circles highlight emission line clumps where this contrary spatial convolution effect is clearly seen.

vary up to $\pm 500 \text{ km s}^{-1}$ from that of the continuum component. The initial continuum fitting was performed using SPS templates from Bruzual & Charlot (2003, hereinafter BC03), with a fixed velocity offset given by the SDSS redshift. We therefore update with the velocity determined from the emission line fitting, and again recompute both continuum and line fitting steps. Due to the discontinuous wavelength coverage and the relatively low spectral resolution, the velocity dispersion of the stellar continuum is difficult to determine, we therefore adopt the velocity dispersion calculated from the SDSS fibre spectrum. Typical values being twice that of the effect the spectral resolution at 5500 \AA . We assume that the stellar velocity dispersion is constant across the whole galaxy, whilst this is not ideal we note that it produces a visibly acceptable result, see Fig. 2. Errors on measured line fluxes are determined by the Levenberg–Marquardt least-squares fitting, however, these formal errors are often an underestimate of the real errors of the line fluxes. Following a procedure derived from SDSS duplicate observations, as discussed by Brinchmann et al. (2013, herein B13), we can translate our formal uncertainties to more representative values. As a result of the aforementioned contrary offsets in the red and blue arms, we modify the standard Gaussian fitting by tying the velocity offsets of the blue and red instrument arms separately.

Spatially binning the 2D spectrum is necessary to optimally extract emission line fluxes. In order to avoid the line broadening effects caused by co-adding spectra with different velocity offsets, we adopt a similar approach to Moran et al. (2010), whereby a two-stage binning strategy is applied. First we adopt a simple binning process, working from the centre of the galaxy outwards, accreting spectra until a minimum continuum S/N of 6 \AA^{-1} is reached. If a bin spans more than 10 arcsec before this threshold is reached, then the binning is terminated. Using the spectral fitting routine, we extract the velocity of the emission lines in the red-arm, where $H\alpha$ is dominant. To this velocity we fit the rotation curve using the following parametrization of Böhm et al. (2004)

$$V(r) = V_{\text{max}} \frac{r}{(r^a + r_0^a)^{1/a}} + V_0, \quad (1)$$

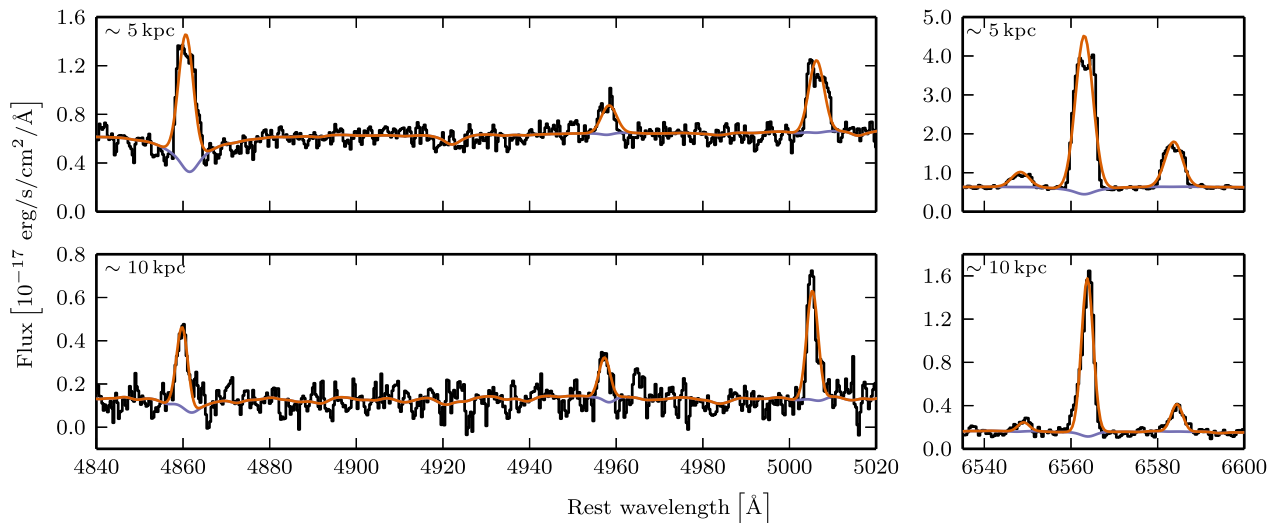


Figure 2. Example of PLATEFIT spectral fitting to emission lines for galaxy BD 5, shown at two radii, $\sim 5 \text{ kpc}$ (top) and $\sim 10 \text{ kpc}$ (bottom). Observed spectrum and its error indicated by black line and shaded area, respectively. The best-fitting model solution is shown in orange, the continuum component is shown in blue. Left: spectra from blue arm covering region containing $H\beta$, $[O \text{ III}]_{\lambda 4959}$ and $[O \text{ III}]_{\lambda 5007}$ lines. Right: spectra from red arm covering region containing $[N \text{ II}]_{\lambda 6548}$, $H\alpha$ and $[N \text{ II}]_{\lambda 6584}$ lines.

where r is the radius, V_{\max} is the maximum velocity at $r \gg r_0$, V_0 is a constant offset velocity, a and r_0 control the shape of the profile. By using a model, we can interpolate the rotation velocity at any position along the spectroscopic slit in a numerically stable fashion. Weighting the velocity measurements by $H\alpha$ S/N, this model provides a reasonable approximation to the true rotation curve, within the limitations imposed by clumpy emission smaller than the slit. Using this velocity fit, the 2D spectrum was shifted to a common rest frame.

Since we are interested in the emission line properties it is ideal to bin spectra on emission line criteria, as opposed to the stellar continuum criteria used previously. Therefore the best-fitting model continuum is first subtracted from the rest-frame shifted spectra, before we apply a custom binning algorithm. Due to the clumpy nature of the emission, any binning algorithm must account for this. With this in mind, we apply a moving boxcar bin to the unbinned spectra, working from the galaxy centre outwards. On each subsequent pass the boxcar is incrementally increased in size. The S/N of each boxcar bin is determined from a least-squares fit of a Gaussian function to the $H\alpha$ line, where a successful bin is when S/N in $H\alpha$ exceeds 6. After the boxcar binning is completed, any remaining unbinned spectra are accreted into the nearest bins, provided their contribution boosts the S/N. To reduce the statistical dependence between neighbouring bins we impose a minimum bin size of ~ 1.6 arcsec, roughly equal to that of the worst FWHM seeing of our observations. The bin centre is defined by the $H\alpha$ -weighted contribution of each 1D spectrum to its respective bin. As the spectroscopic slits were not necessarily aligned with the measured semimajor axes of the galaxies, we additionally deproject these radii, assuming a thin disc and adopting the inclinations and PAs of the galaxies, as derived in Section 2.2.

With the new binning, we apply again our full spectral fitting procedure to the rest-frame shifted spectra. All results in the following are derived from the resulting outputs.

4.1 Bluedisk SDSS properties

In Fig. 3 we compare the central metallicities, central SFRs and total stellar masses of the Bluedisk sample to other galaxies drawn from the 7th data release (DR7) of the SDSS. We discuss the estimation of metallicity in Section 4.3 below, here it suffices to say that the

same estimator has been used for SDSS and the Bluedisk spectra. The Bluedisk sample lies on the same mass–metallicity relation and mass–SFR relation of the DR7 population at a similar redshift. The galaxies have noticeably elevated metallicities with respect to their counterparts of similar SFRs; however, this is to be expected since they are amongst the most massive ~ 10 per cent of galaxies at their epoch. When observed as a function of SFR, a clear split is observed between the $H\text{I}$ -rich and control populations, with $H\text{I}$ -rich galaxies more actively forming stars at their centres. While not an explicit selection criteria of the survey, the two populations do not differ significantly in central metallicity.

4.2 Contamination from non-star-forming sources

To avoid deriving erroneous metallicities, we must take care to exclude spectra contaminated by significant line emission from active galactic nuclei (AGN) or low-ionization nuclear emission-line regions (LINERs).

Following the prescription of Brinchmann et al. (2004), we classify the spectra into five categories. Namely star-forming (SF), AGN/LINER, ‘composite’ SF + AGN/LINER, low S/N AGN/LINER and low S/N SF. For this we use the diagnostic criteria of Kauffmann et al. (2003b) and Kewley et al. (2001), applied to the $([O\text{III}]/H\beta)$, and $([N\text{II}]/H\alpha)$ emission line ratios. However, we find the Kauffmann et al. (2003b) division between SF and composite emission is too aggressive, resulting in the rejection of data points at large radii. Following Moran et al. (2012), we therefore loosen our criteria by offsetting this diagnostic line diagnostic $+0.1$ dex in both $[O\text{III}]/H\beta$ and $[N\text{II}]/H\alpha$ line ratios. Fig. 4 shows the Baldwin, Phillips & Terlevich (1981) (BPT) diagnostic diagram for the co-added spectra. From this we can see that almost all excluded data points are found in the inner regions of the galaxies where we could expect contamination by AGN and shock heated gas.

We ultimately apply an $H\alpha$ S/N ≥ 10 cut to our data. Although we do not apply an equivalent width (EW) cut, we note that 96 per cent of our spectral bins have $EW(H\alpha) > 3$. This EW limit was recommended by Cid Fernandes et al. (2011) to identify pure star-forming galaxies, where $H\alpha$ emission is associated with $H\text{II}$ regions rather than post-asymptotic giant branch stars.

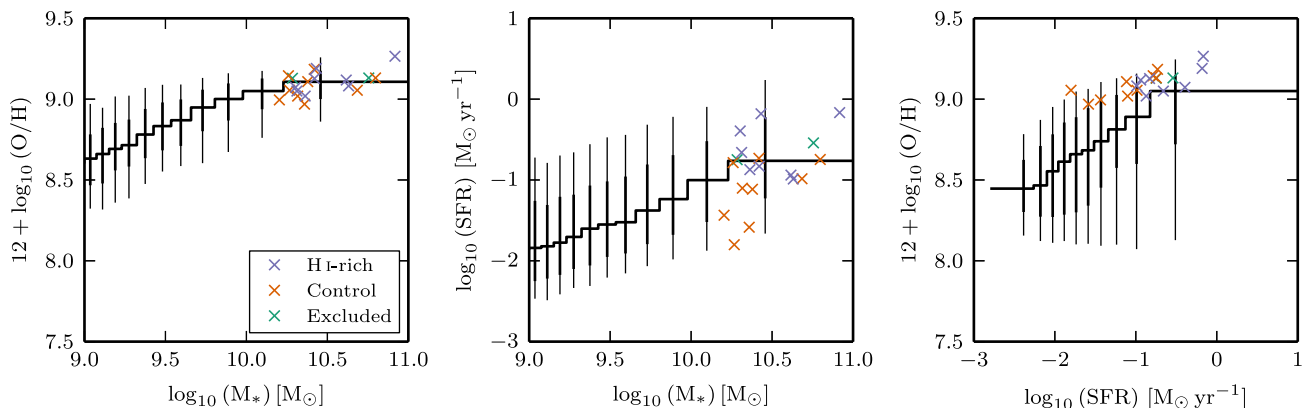


Figure 3. Comparison of Bluedisk galaxy properties with the whole SDSS DR7 population that has all three quantities determined (total galaxy stellar mass, fibre SFR and gas-phase metallicity). The SDSS DR7 sample is selected to span a similar redshift range as the Bluedisk sample ($0.02 \leq z \leq 0.03$). In each plot, crosses indicate SDSS measurements of the Bluedisk $H\text{I}$ -rich, control and excluded samples. In each plot, a stepped-curve represents the median ordinate value of given abscissal bin. Left and middle: sample divided by mass into vigintile bins; right: sample divided by SFR into decile bins. Thick and thin vertical lines represent the 68 per cent (1σ) and 95 per cent (2σ) ranges, respectively.

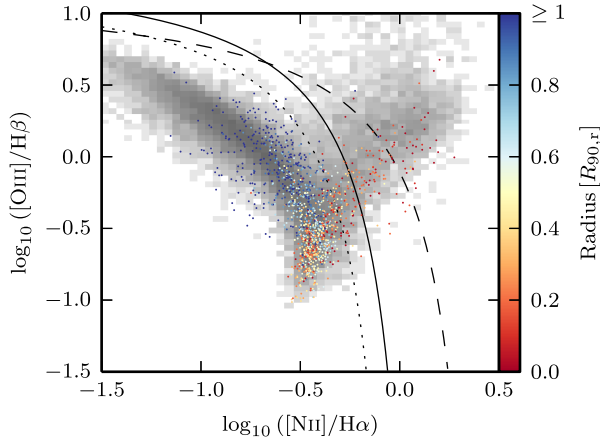


Figure 4. The BPT diagnostic diagram for all binned galaxy data points with $H\alpha$ $S/N \geq 10$ and $S/N \geq 3$ in the other three emission lines. Colours indicate radius from galaxy centre in units of $R_{90,r}$, the radius containing 90 percent of the SDSS r -band luminosity. Shaded underlay shows BPT histogram of SDSS DR7 fibres in the same redshift range ($0.02 \leq z \leq 0.03$). The dashed line indicates the division between pure AGN and ‘composite’ galaxies (Kewley et al. 2001). The dotted line represents the nominal Kauffmann et al. (2003b) SF diagnostic line, but for our classification we use a modified form of this offset +0.1 dex in both axes (solid line).

4.3 Inferring metallicities

We derive gas-phase metallicities using the method developed by B13. This method applies a Bayesian framework to a grid of photoionization models. We shall pair the B13 methodology with the photoionization models of Charlot & Longhetti (2001, herein CL01). To correct for dust, the B13 procedure uses a two-component dust-absorption model of Charlot & Fall (2000), with a wavelength-dependant attenuation curve of the form $\tau(\lambda) \propto \lambda^{-1.3}$. For a set of emission line fluxes,³ we obtain a posterior probability distribution on the metallicity. However, to test our CL01 derived metallicities we will now outline a series of other metallicity determination methods.

A common set of methods for inferring metallicities is to use line-ratio diagnostics that have directly calibrated to oxygen abundances of $H\ II$ regions (either theoretical or observed). We use two such methods, one using the theoretically derived relations of Kobulnicky & Kewley (2004, herein KK04), and another method using the empirically derived relations of Pilyugin & Mattsson (2011) known as the NS calibrator. In addition, we check the sensitivity of our CL01 metallicities to the choice of photoionization models, by applying the models of Dopita et al. (2013, herein D13) within the B13 framework. It should be noted that the B13 method is similar in principle to others such as IZI (Blanc et al. 2015) and HII-CHI-MISTRY (Pérez-Montero 2014).

Dust attenuation affects line ratios, therefore we must correct for dust before we apply the KK04 and NS calibrators. We adopt the same $\tau(\lambda) \propto \lambda^{-1.3}$ attenuation curve and calculate its normalization by assuming an intrinsic $H\alpha/H\beta$ Case-B ratio of 2.85 (temperature, $T = 10^4$ K, and electron number density, $n_e = 10^4$ cm⁻³) (Osterbrock & Ferland 2006). We note, however, that by assuming such physical properties of the $H\ II$ regions we must a priori assume a metallicity. With the B13 approach, we avoid this assumption by simultaneously correcting for dust attenuation when

inferring metallicity. Nevertheless, for spectra with high S/N , both the B13 and empirical Case-B approaches yield similar results for the strength of the attenuation.

Uncertainties on the metallicity determinations for the KK04 and NS methods are determined using Monte Carlo simulations. We assume the true line fluxes to be normally distributed about the measured line flux, with a standard deviation equal to the error in the measured value. Of these many realizations, we take the median as the metallicity value, and the symmetrized $\pm 1\sigma$ quantiles to be its associated error. For the CL01 and D13 models, we can extract the median and its error directly from the cumulative posterior probability of the metallicity parameter. Due to the finite sampling of the metallicity parameter in these models, we impose an additional minimum uncertainty of ± 0.05 dex (Brinchmann et al. 2004). We additionally apply the same minimum error to the KK04 and NS methods.

In Appendix A, we compare the different metallicity methods. We show that the CL01 method produces results consistent with the other three methods. We adopt CL01 as our default method for metallicity determination. Herein for simplicity when we refer to metallicity, we are referring to that which is derived from the CL01 models.

Note that in contrast to some studies of abundance gradients our spatial bins include line emission from both $H\ II$ regions and diffuse emission. Since previous studies have found that the diffuse emission is powered by radiation escaping from $H\ II$ regions (e.g. Hoopes & Walterbos 2003; Kim et al. 2013), we can approximately treat the combined line emission as coming from an $H\ II$ region with larger volume and hence a lower ionization parameter. The CL01 models are well suited for this as they cover a range in ionization parameters. In addition, note that while the CL01 models are not specifically optimised for spatially resolved regions in galaxies, tests in B13 showed that the CL01 models perform well in this case.

4.4 Estimating gas mass densities

To determine gas surface mass density at the same resolution as our metallicities, we estimate the gas surface mass densities directly from our spectra. In B13, it was shown that when most of the strong lines in the optical spectrum are available it is possible to use photoionization models with a flexible treatment of metal depletion to place constraints on the gas surface mass density of galaxies. The application shown in B13 used the CL01 models which we also use here. By jointly fitting the strong optical lines B13 showed that the total gas surface mass densities can be estimated through

$$\Sigma_{\text{gas}} = 0.2 \frac{\tau_V}{\xi Z} M_{\odot} \text{pc}^{-2}, \quad (2)$$

where τ_V is the optical depth in the V band, ξ the dust-to-metal ratio of the ionized gas, and Z the metallicity.

They compare the result of applying this relation to spectra from the SDSS to total mass densities measured from $H\ I$ and H_2 mass maps from the THINGS (Walter et al. 2008) and HERACLES (Leroy et al. 2009) surveys. This point-by-point comparison showed that the spectroscopic method is in excellent agreement with the $H\ I+H_2$ mass maps, except at the very highest gas surface densities, $\Sigma_{\text{gas}} > 75 M_{\odot} \text{pc}^{-2}$.

For this paper, we note that we use the same set of emission lines used by B13 in their study with comparable S/N , so we expect this result to carry over to our study. This means that our spectroscopic gas densities are likely to be underestimated in the central regions

³ The B13 analysis was performed using emission lines: $[O\ II]_{\lambda 3727}$, $H\delta$, $H\gamma$, $H\beta$, $[O\ III]_{\lambda 4959}$, $[O\ III]_{\lambda 5007}$, $H\alpha$, $[N\ II]_{\lambda 6584}$, $[S\ II]_{\lambda 6716}$ and $[S\ II]_{\lambda 6731}$.

of the galaxies. We expand on this and discuss this method more in Appendix B.

5 RESULTS

5.1 The local mass–metallicity relation and radial mass profiles

From basic analytical arguments one expects the metallicity of a system to depend on the stellar and gas mass budgets (Pagel 1997). Indeed, recent works have uncovered a correlation between stellar mass surface density and metallicity, known as the local mass–metallicity relation (Moran et al. 2012; Rosales-Ortega et al. 2012; Sánchez et al. 2013). However, it is worth noting Pilyugin et al. (2014) find that although local surface brightness and metallicity are correlated, there is no unique relation between the two that holds at all radii in galaxies. Nevertheless, we will now test whether the local mass–metallicity relation holds for the Bluedisk galaxies.

In Fig. 5, we present radial profiles of our spectroscopic dust-to-gas estimates of gas-mass surface densities, Σ_{gas} , as well as our stellar mass surface densities, Σ_* , which are matched in aperture. The median trends for Σ_* exhibit clear differences between the H I-rich and control samples, with H I-rich being consistently less

massive at a given scale radius. By contrast, Σ_{gas} shows no significant distinction between the samples. We caution that although there appears to be a slight upward trend in Σ_{gas} with radius, which may be artificial (see Section 4.4).

In Fig. 5, we also show the radial profiles of the gas-to-stellar mass ratio, $r_{\text{gas}} = \Sigma_{\text{gas}}/\Sigma_*$. Here, we find that the H I-rich galaxies exhibit enhanced r_{gas} ratios at all but the very centre the stellar disc. We note that this is primarily driven by radial differences in Σ_* rather than Σ_{gas} . In other words, at a fixed Σ_* the H I-rich and control samples are indistinguishable in terms of r_{gas} .

In Fig. 6, we show the local mass–metallicity relation for the Bluedisk galaxies. We also plot the correlation between r_{gas} and metallicity. Crucially, neither of these correlations show any strong offsets between the H I-rich and control samples, implying that the processes that govern these parameters are similar in both classes of galaxies. We observe that the correlation of metallicity with r_{gas} is not visibly tighter than that with Σ_* . But, as exemplified in Fig. 6(c), we note that at the lowest stellar mass densities ($\log_{10}(\Sigma_*) \lesssim 1.5 \text{ M}_{\odot} \text{ pc}^{-2}$) a significant portion of the metallicity variation can be explained by changes in Σ_{gas} . We must caution, however, that metallicity and Σ_{gas} are not independently derived. The nature of the modelling will introduce a small intrinsic

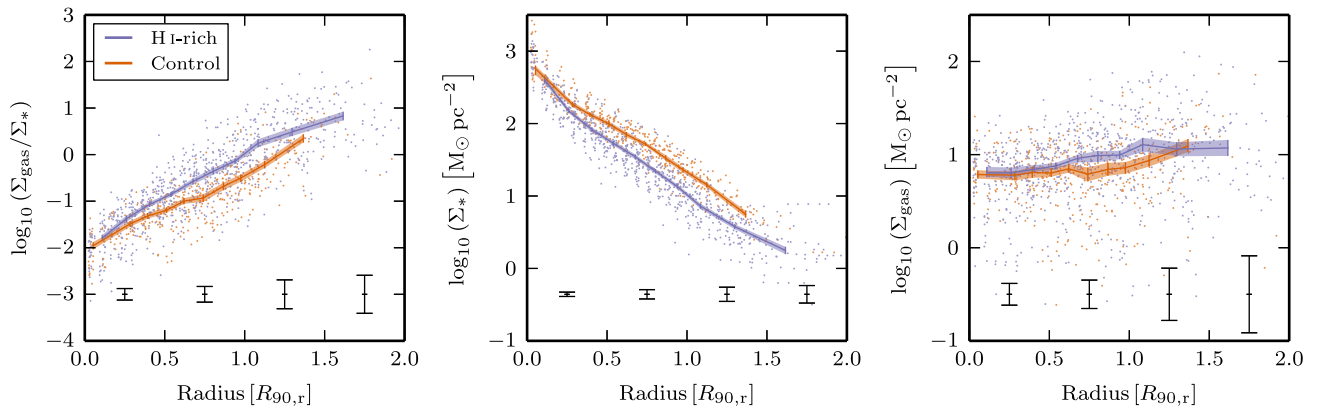


Figure 5. Radial profiles of gas-to-stellar mass ratio, $r_{\text{gas}} = \Sigma_{\text{gas}}/\Sigma_*$, stellar mass density, Σ_* , and gas mass density, Σ_{gas} , are shown from left to right. Data are shown for both Bluedisk H I-rich and control samples, coloured blue and orange, respectively. We plot the individual data points, as well as the binned median trends in each plot. Shaded regions indicate $\pm 1\sigma$ errors on the trends, as determined by bootstrapped Monte Carlo realizations. Radius is in scale units of $R_{90,r}$, the radius containing 90 per cent of the SDSS r -band luminosity. Black vertical bars indicate median error in each $0.5 R_{90,r}$ division.

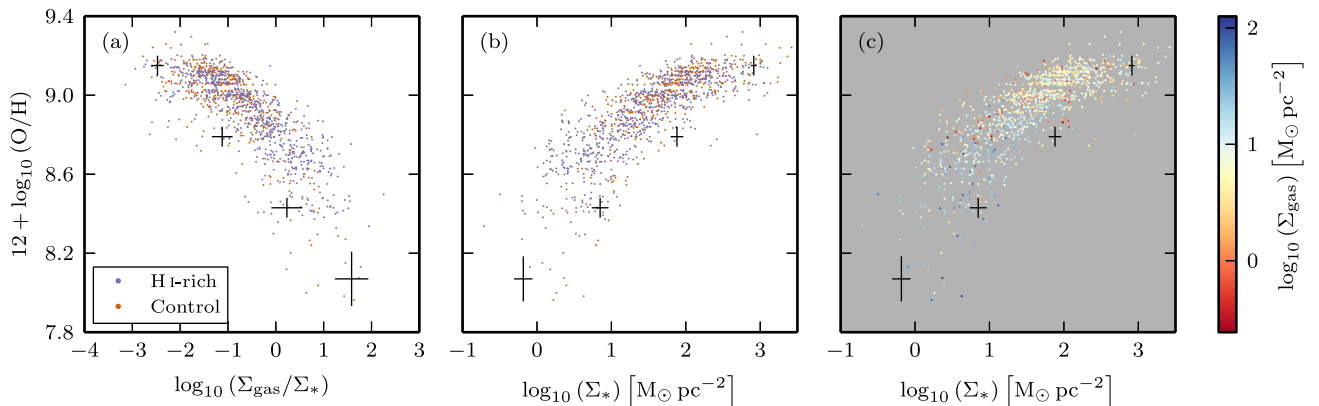


Figure 6. Correlations between metallicity and surface mass densities. (a) Metallicity against gas-to-stellar mass ratio. (b),(c) both show metallicity against stellar surface mass density Σ_* , i.e. the local mass–metallicity relation, but are colour-coded in differently. In (a), (b) we colour according to Bluedisk sample, H I-rich and control samples, distinguished by blue and orange data points, respectively. In (c), we colour the local mass–metallicity relation by gas surface mass density, Σ_{gas} . Black crosses indicate median error within equally spaced bins.

correlation between these two parameters. The magnitude of this effect is not easily quantified.

Having shown that both H I-rich and control samples form a consistent local mass–metallicity relation, we shall explore the radial metallicity profiles of the Bluedisk galaxies.

5.2 Metallicity profiles of the Bluedisk galaxies

We present the metallicity profiles of the Bluedisk galaxies in Figs 7, 8 and 9, divided into their H I-rich, control and excluded samples, respectively.⁴ By visual inspection alone there is no clear distinction between H I-rich and control samples in terms of the profile shapes. The control sample does appear to be more radially truncated, however, this is to be expected. It has been shown by Wang et al. (2011) that H I mass fraction is correlated with the $g - i$ colour gradient in galaxies. This implies that the H I-rich galaxies have higher rates of star formation in their outer discs. Therefore we expect the control sample to have less extended star formation, ultimately limiting the radius out to which we can robustly detect emission lines. For a more quantitative analysis, we must study the measured gradients of the linear model.

As is common in the literature, we approximate the metallicity profiles with a symmetric linear best fit. In most cases, this simple functional form encapsulates the overall change in metallicity from the centre to the outskirts of the galaxy. In Fig. 10, we show the dependence of the metallicity gradient on both stellar mass and the H I mass fraction. There exists a significant correlation between H I mass fraction and metallicity gradient. Galaxies with larger H I mass fractions typically have steeper metallicity gradients. Whereas we find no significant correlation between stellar mass and metallicity gradient. However, this is unsurprising given the narrow range that we span ($10.2 \lesssim \log_{10}(M_*/M_\odot) \lesssim 11.0$). We note these results remain unchanged when we adopt a different scale radius, R_{25_g} (the radius at which the SDSS g -band surface brightness reaches $25 \text{ mag arcsec}^{-2}$). We highlight two galaxies with especially steep metallicity gradients. First, BD 34 which shows very large errors in its measured metallicity gradient. This gradient is measured from only four valid data points and is poorly constrained. Secondly, BD 45 whose metallicity profile shows some hints of asymmetry, but otherwise offers no explanation for the excessively steep inferred metallicity gradient. Regardless, we retain both these outlying galaxies in our analysis.

From visual inspection, describing some of these galaxies with a straight-line model appears to poorly reflect the true metallicity profile. In a number of galaxies the metallicity gradient appears to be increasing with radius. As such, a gradient measured from the outer disc would be much steeper than one measured from the inner disc. With long-slit spectra, we only measure metallicity along one dimension of the galaxy. So if significant azimuthal metallicity variations are present, our metallicity measurements may not be indicative the whole galaxy at a given radius. Simulations of Petit et al. (2015) indicate, however, that azimuthal variations decay are expected to decay on time-scales shorter than the orbital period of the galaxy. Indeed, observationally there is little support for strong azimuthal variations, with Sánchez et al. (2015) reporting only modest ($<0.05 \text{ dex}$) azimuthal variations.

In the following, by stacking the metallicity profiles we shall attempt to produce average metallicity profiles.

5.2.1 Stacked average metallicity profiles

To study the metallicity profiles for the Bluedisk galaxies further, we stack the individual metallicity data points into equal mass decile radial bins. Tracing the median metallicity of the bins we construct the average metallicity profiles, which are shown in Fig. 11. We caution that since the metallicity data points are equally weighted, the outermost bin of each stack might be considered unreliable (see Appendix C).

In Fig. 11(b), we show the stacked profiles of the H I-rich and control samples. We also bisect each sample by total stellar mass. We observe that all galaxies have similar central metallicities, but different profile shapes. We note that the outermost bin of the high-mass control profile has a spuriously low metallicity and should be ignored (see Appendix C). Putting this aside, the stacked profiles appear to indicate a shallower inner gradient and a steeper outer gradient. It is difficult to define the characteristic radius at which this transition happens. However, by eye it seems that the transition occurs at a smaller radius in the H I-rich galaxies than in the control galaxies. Overall, we note that the transitions in the stacked profiles do not appear as abrupt as they do in the unstacked profiles. This would imply that using $R_{90,r}$ as a radial coordinate is not ideal for expressing this turnover. Indeed, since we observe a local mass–metallicity relation, a scale radius based on stellar mass density would perhaps be more appropriate.

Beyond H I characteristics there are other aspects which may affect metallicity profiles. Using semi-analytical models Fu et al. (2013) predict metallicity gradients should be correlated with B/T ratio. In these models, galaxies with more prominent bulges are expected to have shallower gas phase metallicity gradients as the gas distribution in these galaxies is set by later infall of gas. We show the stacked profiles of the galaxy samples bisected by B/T light ratio in Fig. 11(c). We observe no apparent connection between metallicity profiles and bulge prominence, however, given the low bisecting threshold ($B/T = 0.15$), we find our results to only be in mild tension with these predictions. These observations are consistent with Sánchez et al. (2014), who observe no correlation between metallicity gradient and morphological galaxy type.

It may be possible that the differences we observe between our H I-rich and control populations are driven by the effects of bars. It has been established that there exists an anticorrelation between H I mass fraction and the bar occurrence fraction (Davoust & Contini 2004; Masters et al. 2012). Numerical simulations have also shown that the presence of galaxy bars can drive enhance gas mixing, flattening the metallicity profile (Friedli, Benz & Kennicutt 1994; Minchev et al. 2011). These effects of bars on metallicity been borne out by observations (e.g. Martin & Roy 1994; Dutil & Roy 1999). However, it has also been shown that when metallicity gradients are expressed units of effective disc radius, rather than physical distance, bars show no significant impact on the metallicity gradient (Sánchez et al. 2014). Even so, it is prudent when comparing the metallicity profiles of the H I-rich and control populations that we take care to exclude the potential impact of bars.

By visual inspection, we classify 30 percent of the Bluedisk galaxies to be strongly barred (with four galaxies indeterminate due to high inclination). This rate is consistent with galaxies of the same stellar mass (Skibba et al. 2012). In Fig. 11(d), we show stacked metallicity profiles of the unbarred galaxies. The distinction between H I-rich and control samples clearly remains after excluding barred galaxies. In Fig. 11(f), we show the effects of bars on the whole sample, and in Fig. 11(g) the effects of bars exclusively

⁴ Larger versions of the metallicity profiles in conjunction their SDSS gri composite images are available in Appendix D (online only).

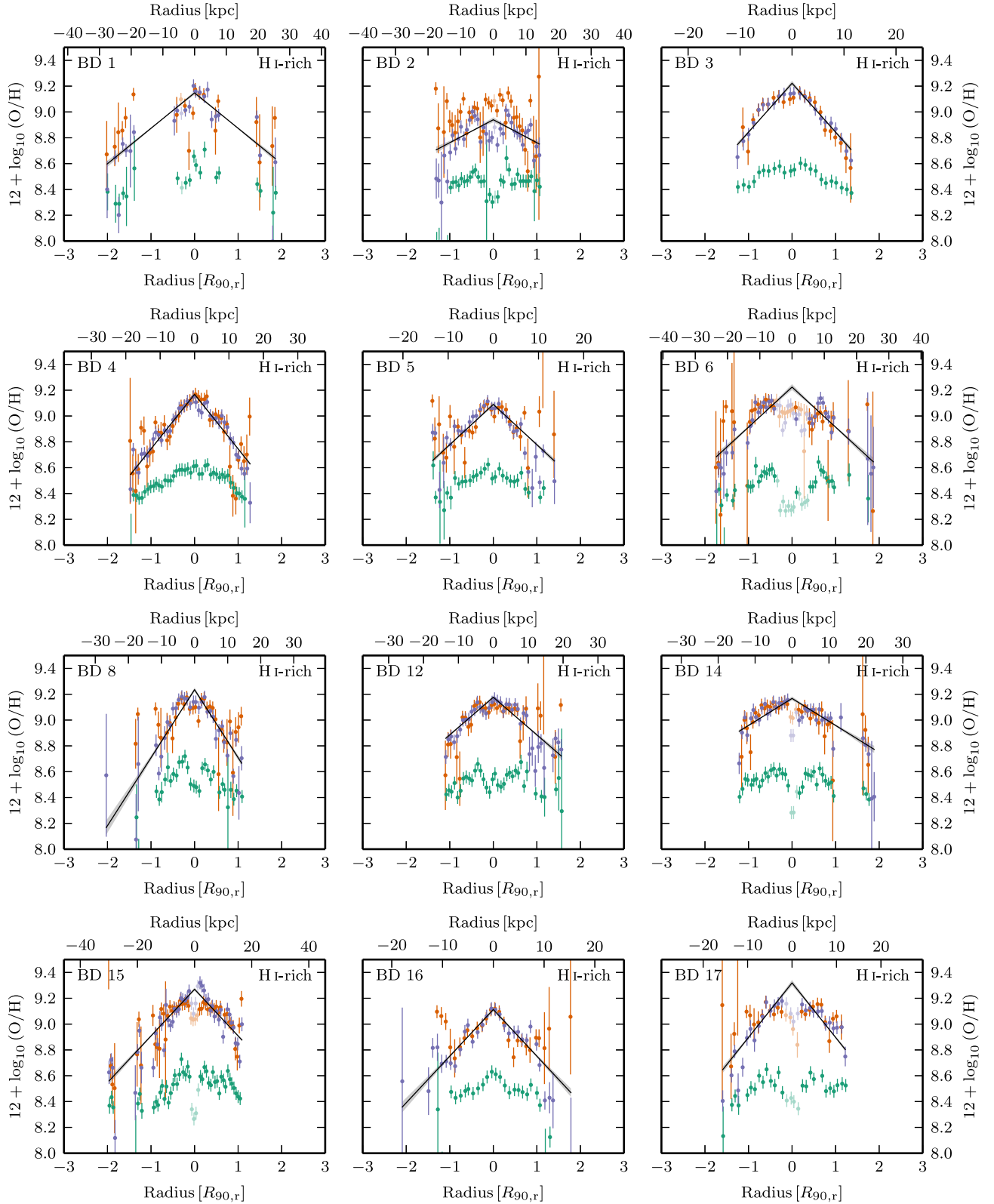


Figure 7. Metallicity profiles of the 23 galaxies in the Bluedisk H I-rich sample. The metallicity is inferred using the CL01, KK04 and NS methods, plotted as blue, orange and green, respectively. For the CL01 metallicities, we show the best-fitting straight-line model (black), where the shaded area indicates its associated $\pm 1\sigma$ error in gradient. We also show in pale colours data points masked from the fitting due to their non-SF emission characteristics. All plotted data has an minimum $H\alpha$ S/N ≥ 10 . Radius is plotted in both units of a scale radius, $R_{90,r}$, and physical size in kpc.

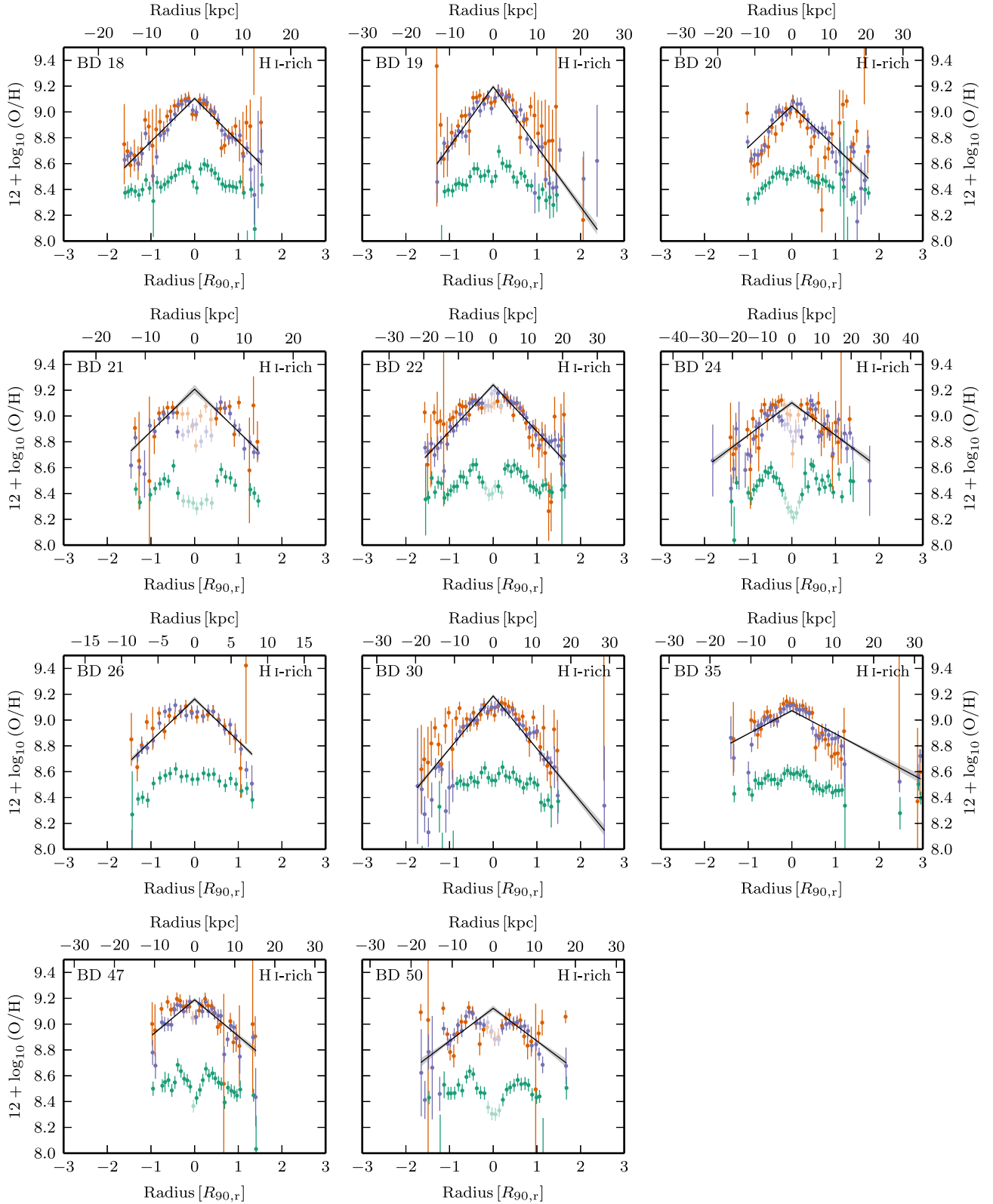


Figure 7 – continued.

on the control sample. From both of these figures, we notice some flattening of the metallicity profile in galaxies with bars. We caution the reader that our sample of H I-rich galaxies with bars is very small.

We have repeated the stacking analyses for the other metallicity calibrators. We see similar effects when using the KK04 method, but we do not observe the outer metallicity drops when using the NS method. It should be noted, however, that the NS method does

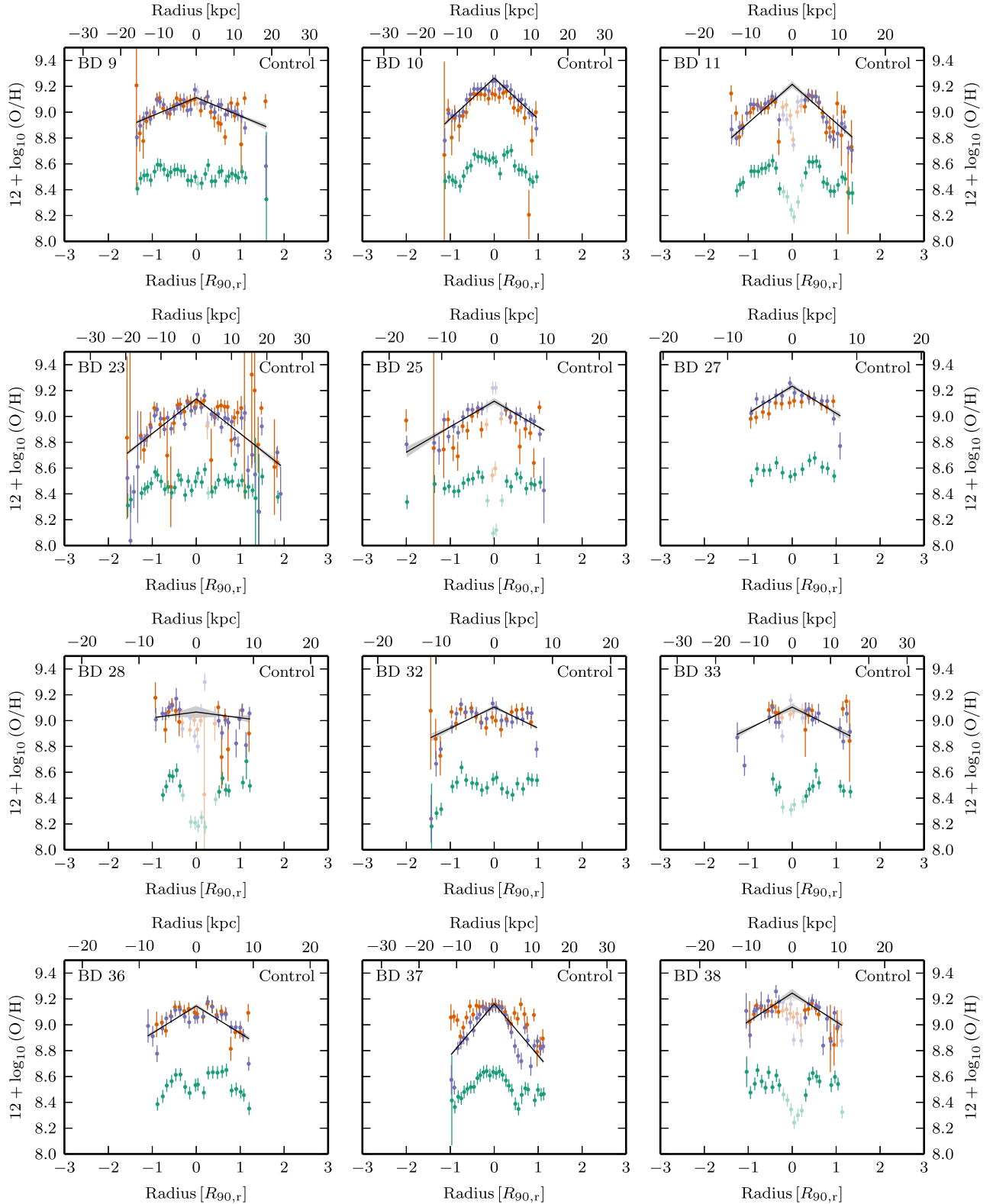


Figure 8. Metallicity profiles of the 18 galaxies in the Bluedisk control sample, see Fig. 7 for details.

not allow for variations in the ionization parameter. Methods that do not include this extra dimensionality may not be best suited for the study we present here.

Finally, we note in Fig. 11(a) that there appears to be a significant amount of scatter in the inner regions of the galaxies. We identify

galaxies that harbour AGN using the $(\text{O III}/\text{H}\beta)$ and $(\text{N II}/\text{H}\alpha)$ emission line ratios from the centre of each galaxy. We adopt the criterion of Kauffmann et al. (2003b) and identify galaxies with central non-SF emission. We find that galaxies with central non-SF emission exhibit a two-fold increase in the scatter of the metallicities in the

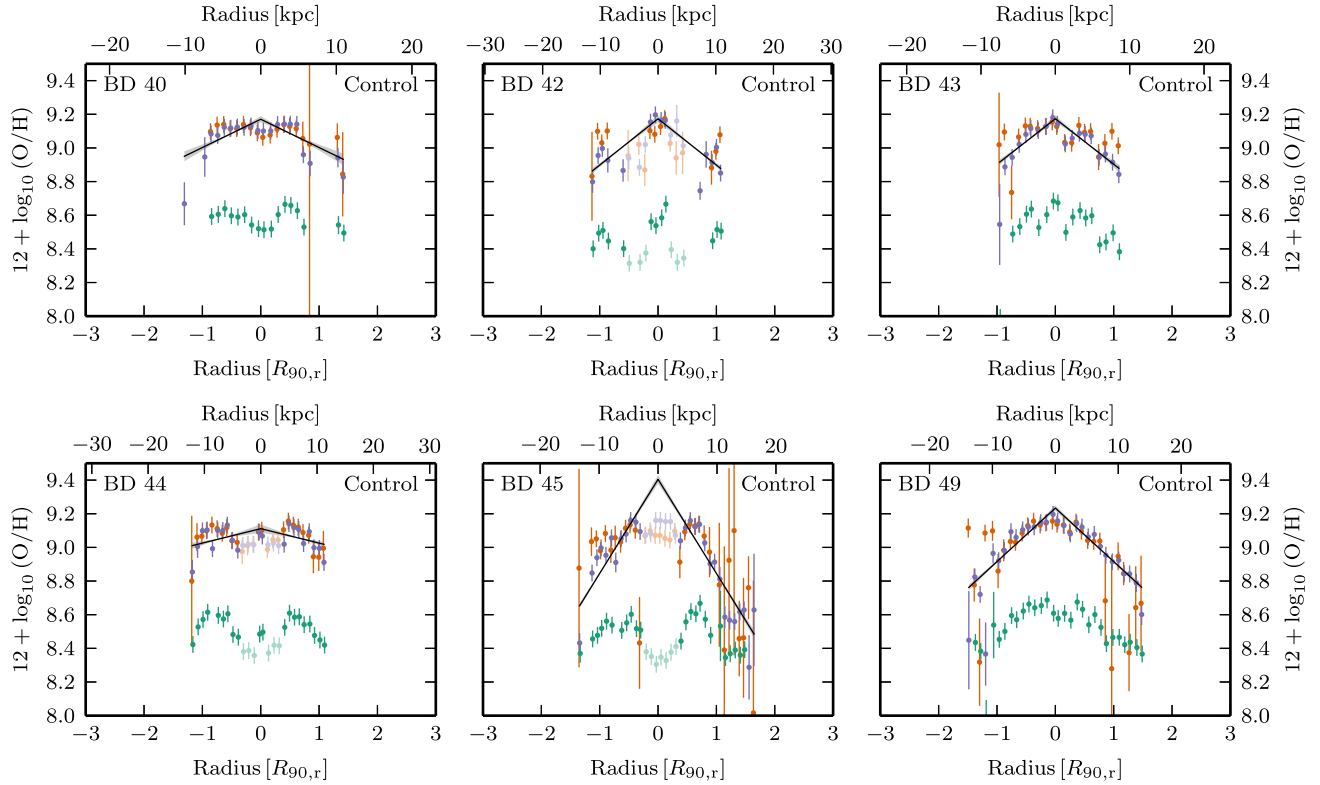


Figure 8 – continued.

inner region ($r < 0.3 R_{90,r}$). This might be an indication of AGN interacting with the central environment. However, it is equally plausible that we are not sufficiently excluding non-SF contaminated data points, producing erroneous metallicity estimates. Our long-slit spectroscopic observations are not ideal for such study of metallicity scatter, integral field spectroscopy with good spatial resolution may provide sufficient data to study both the radial and azimuthal metallicity scatter in the inner regions of galaxies.

5.3 Summary of results

In the next section, we shall construct a simple analytical model to explain the metallicity profiles we have observed. But first, we shall briefly summarize our results as follows.

- (i) We reproduce the recently reported local mass–metallicity relation. However, we highlight that at low stellar mass densities, there appears to be a residual correlation of metallicity with gas-mass density.
- (ii) The metallicity gradients of a galaxy is strongly correlated with its $H\text{I}$ mass fraction.
- (iii) We stack the metallicity of the galaxies and derive average profile shapes. We find different average profiles for the $H\text{I}$ -rich and control galaxy samples.
- (iv) We find galaxies in both samples that exhibit transitions from shallower inner metallicity gradients to steeper outer metallicity gradients.
- (v) Barred galaxies appear to have flatter metallicity profiles, but this effect does not drive the difference observed between the Bluedisk samples.

6 DISCUSSION

Up to this point we have mainly concerned ourselves with the similarities and differences between our Bluedisk samples. We have, however, not yet suitably tackled the complex issue regarding the origin of the metallicity profile itself. Exploring quantitatively the interplay of the many potential mechanisms is challenging. However, recent years have seen the emergence of a class of simple ‘reservoir’ models (Bouché et al. 2010) in which stars form from a gas reservoir regulated by the star formation and gas flows in and out of the system. While simple, these models are able to provide simple descriptions of the (central) metal content of galaxies at low redshift (Finlator & Davé 2008; Davé, Finlator & Oppenheimer 2012; Lilly et al. 2013). The models are also naturally interpreted as a result of galaxies being generally close to equilibrium between star formation, outflows and inflows (Davé et al. 2012).

6.1 Modelling resolved galaxies as local gas regulators

The reservoir models generally consider the galaxies to be spatially unresolved. Here, we therefore will develop a simple extension of these models to a resolved galaxy. In particular, we will extend the ‘gas regulator’ model (Lilly et al. 2013, herein L13) which has been shown to successfully fit the central metallicities of star-forming galaxies in the SDSS. Our approach will be to minimally extend this model to 2D to see whether such a simple extension is sufficient to describe the metallicity profiles of our galaxies.

To do this, we envisage our disc galaxy divided into a set of radial zones. We then assume that the mean properties of each radial zone can be described by individual gas reservoir models. Alternatively one might take this to mean that we assume that each radial zone

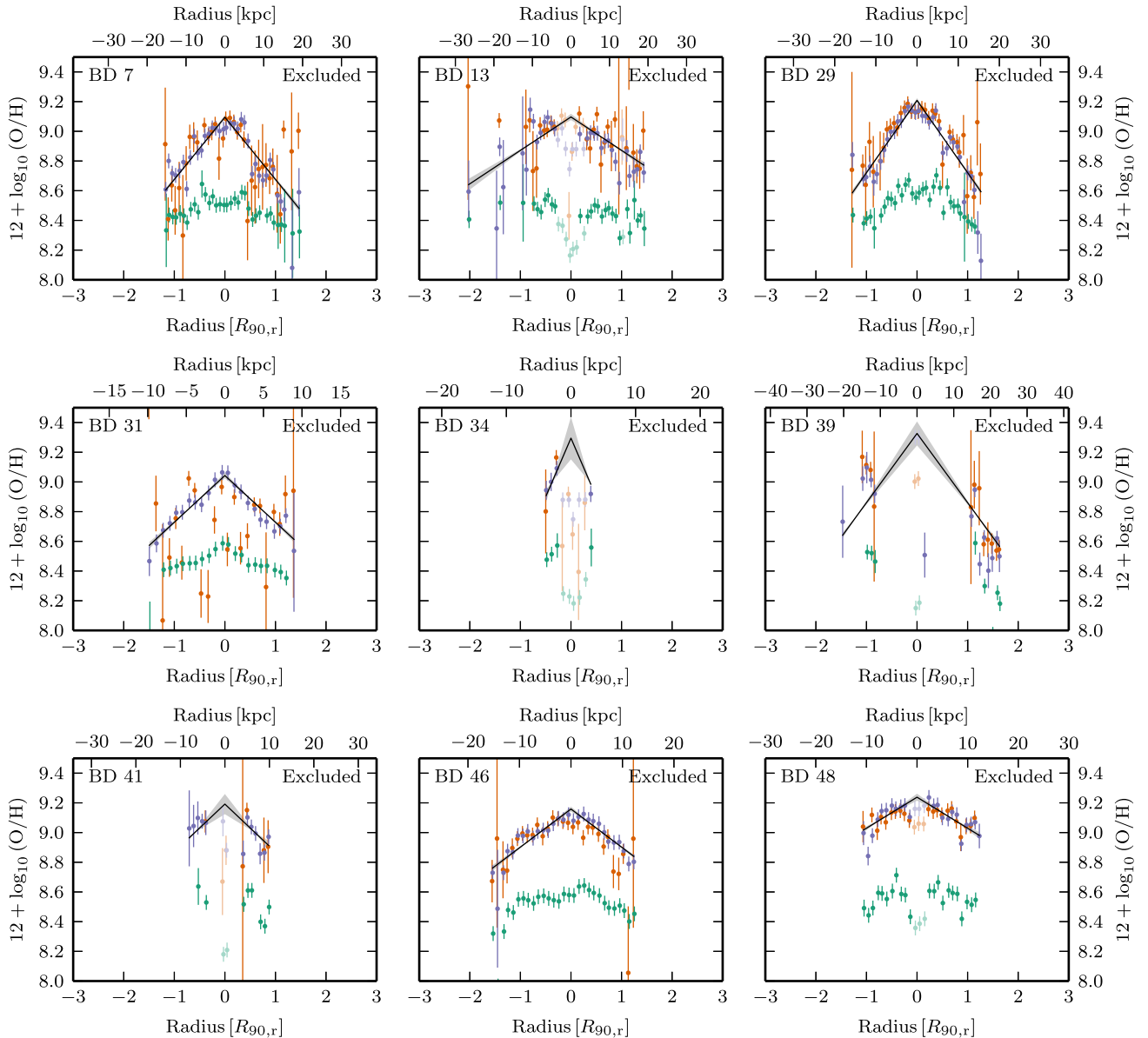


Figure 9. Metallicity profiles of the nine galaxies in the excluded Bluedisk sample, see Fig. 7 for details.

individually is in an equilibrium between inflow, outflow and star formation – a detailed balance principle which is not required by the reservoir models in general. Note that we do not assume that such an equilibrium holds at each point but rather in an average sense across a radial bin.

Generally there might be radial mass transfer between these zones, but in the following we will make the simplifying assumption that radial mass transfer can be ignored. Since, semi-analytic models of Fu et al. (2013) have argued that gas flows are of minor importance. It should be noted, however, that the simulations of Minchev et al. (2011) have shown that in the presence of bars, gas can be efficiently transported resulting in flattened metallicity gradients.

In addition to gas flows, long-lived stars are expected to migrate from their original radius, particularly in the presence of bars (Roškar et al. 2008; Di Matteo et al. 2013). However, in the model we will assume that the mass of stars observed at a given radius represents the total mass of stars formed that given radius. Or rather,

we assume the stars remain associated with the gas from which they form.

We now outline how we adapt and apply the gas regulator model to our data. We refer the reader to L13 for a full treatment and derivation of the model.

6.1.1 Transport of gas

The underlying equation describing the rate of change of the reservoir gas mass (in each radial bin) is

$$\dot{m}_{\text{gas}} = \dot{m}_{\text{in}} - \dot{m}_{\text{out}} - \dot{m}_{*} + \dot{m}_{\text{return}} + \dot{m}_{\text{radial}}. \quad (3)$$

The components are as follows.

(i) \dot{m}_{in} is the rate of metal-poor gas inflowing from the halo to the reservoir. We do not explicitly parametrize \dot{m}_{in} and it shall be eliminated in due course.

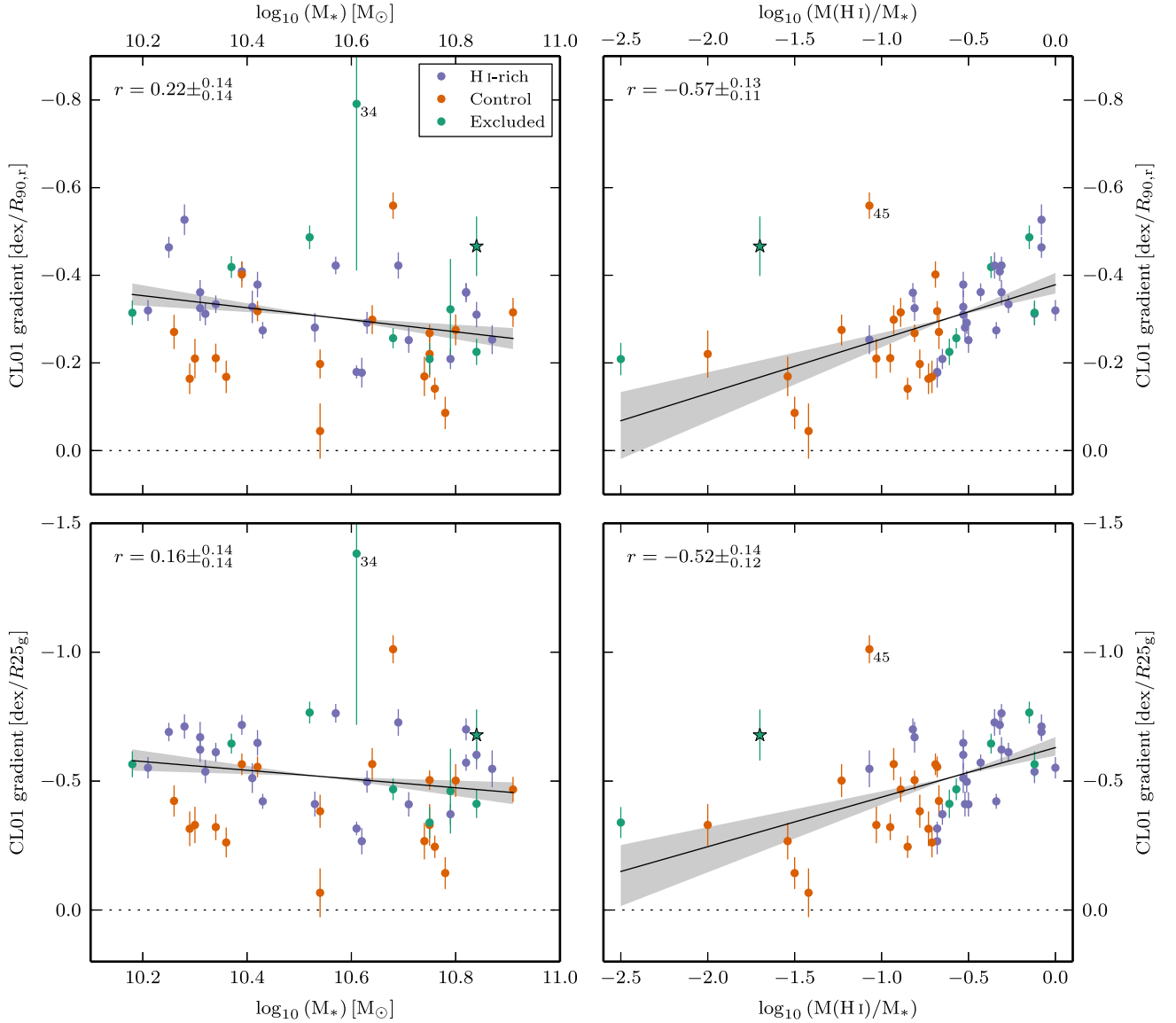


Figure 10. CL01 metallicity gradients as a function of total stellar mass (left) and H I mass fraction (right). We report gradients using two different scale radii, $R_{90,r}$ (top) and $R_{25,g}$ (bottom). Colours indicate Bluedisk sample classification H I-rich (blue), control (orange) and excluded (green). Best linear fit and its error in slope are indicated by the solid lines and the shaded regions, respectively. The Spearman's rank correlation coefficient is given in the top-left corner of each figure. A star indicates BD 39, which is excluded from the regression and the r -statistic computation. Numbers label individual galaxies with especially steep metallicity gradients that we reference in the text. For reference, $r = 0.29$ is the two-tailed Spearman's r -value at a $\alpha = 0.05$ significance level.

(ii) \dot{m}_{out} is the rate at which gas flows out from the reservoir and into the halo or beyond. Since we are mostly concerned with the star-forming disc, this is assumed to be driven by winds from massive stars, and we therefore consider \dot{m}_{out} to be linearly proportional linearly proportional SFR, i.e. $\dot{m}_{\text{out}} = \lambda \text{ SFR}$, where λ is the mass-loading factor. In Section 6.1.3, we attempt to estimate this mass-loading factor.

(iii) \dot{m}_* is the rate at which gas is converted into stars. In other words $\dot{m}_* = \text{SFR}$. The SFR is itself assumed to be linearly proportional to the current mass of the reservoir, $\text{SFR} = \epsilon \dot{m}_{\text{gas}}$, where ϵ is the star formation efficiency. This link between \dot{m}_{gas} and SFR provides the regulatory aspect of the model.

(iv) \dot{m}_{return} is the rate at which enriched gas is returned from short-lived high-mass stars. A fraction R of the mass converted into stars is assumed to be instantaneously recycled back into the reservoir ($\dot{m}_{\text{return}} = R \text{ SFR}$).

Following L13, we adopt a fixed value of the return fraction $R = 0.4$. As shown by BC03 this is the mid-range value over a variety of initial mass functions (IMF), for a 10 Gyr stellar population. Provided there are no strong age gradients across the galaxies it is reasonable to adopt a radially constant return fraction. The exact value of the return fraction will depend on the choice of IMF. We note, however, that our conclusions are not sensitive the exact value we adopt for R . This insensitivity results from the degeneracy of R with parameters that we shall fit (see Section 6.1.4).

(v) \dot{m}_{radial} is the rate at which radial flows within the disc change the gas content of the reservoir. For simplicity we assume $\dot{m}_{\text{radial}} = 0$, neglecting the effects of radial flows.

With these principle assumptions, equation (3) can be written as

$$\dot{m}_{\text{gas}} = \dot{m}_{\text{in}} - (1 - R + \lambda) \text{SFR}. \quad (4)$$

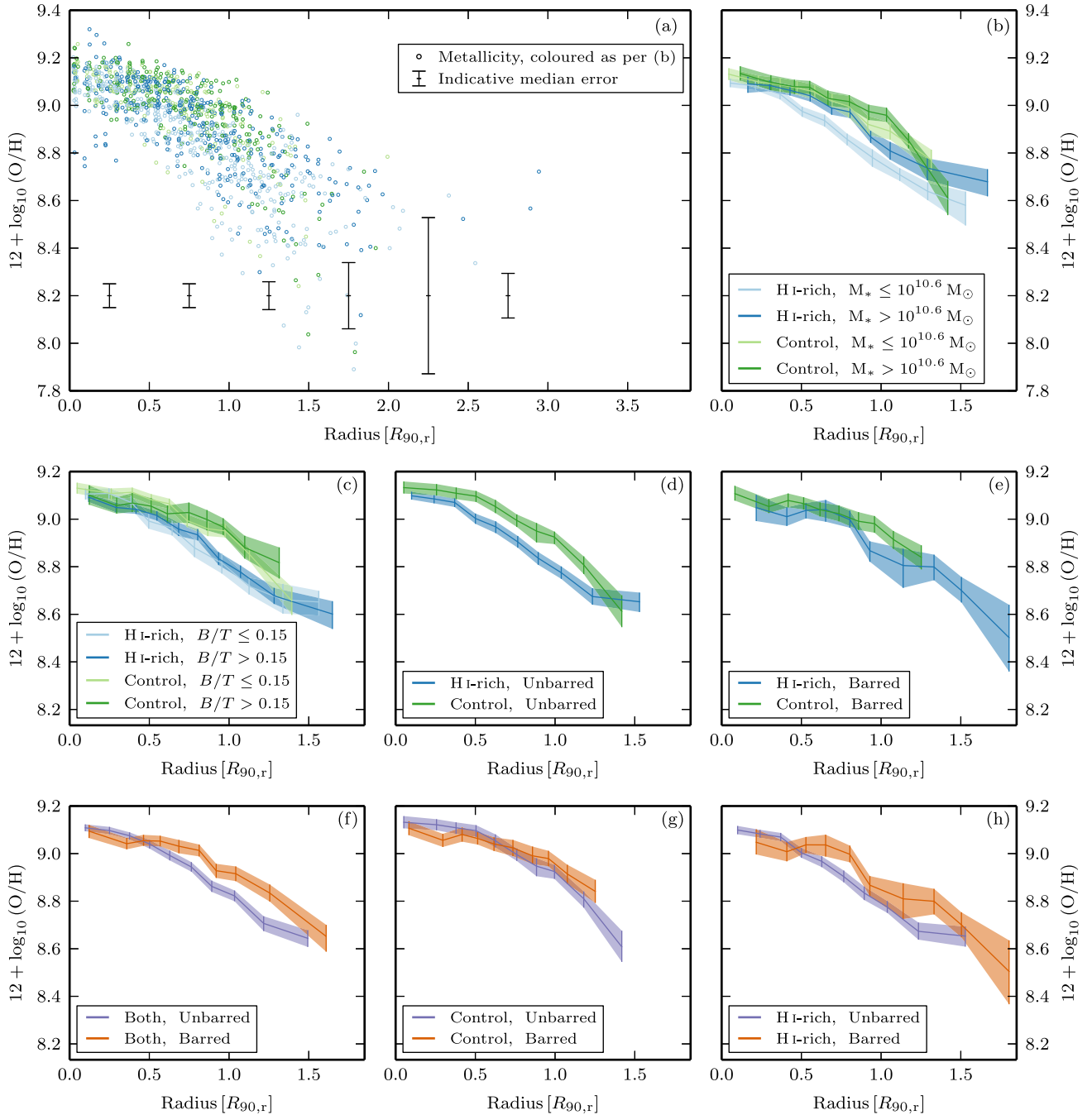


Figure 11. Results from stacking metallicity profiles. (a), (b) Metallicity of the Bluedisk galaxies, colour distinguishing H I-rich and control sample, as well as low and high stellar mass. (a) The individual metallicity data points. Vertical bars indicate median error in each $0.5 R_{90,r}$ division. (b) Median value of galaxies stacked in decile radial bins. Shaded regions indicate $\pm 1\sigma$ errors determined by bootstrapping Monte Carlo realizations. (c) Similarly stacked profiles, however, split by B/T ratio instead of stellar mass. Dividing the H I-rich and control galaxies into barred and unbarred samples we generate the panels (d)–(h).

Furthermore L13 show that by introducing the variable $r_{\text{gas}} = m_{\text{gas}}/m_*$, the ratio of gas-to-stellar mass, equation (4) can be conveniently rewritten as

$$\dot{m}_{\text{in}} = \left((1 - R)(1 + r_{\text{gas}}) + \lambda + \epsilon^{-1} \frac{d \ln(r_{\text{gas}})}{dt} \right) \text{SFR}, \quad (5)$$

which makes the regulatory link between the SFR and gas inflow explicit. We will assume that this holds in the mean in each radial bin.

6.1.2 Transport of metals

We now consider the flow of metals into and out from the reservoir. In the absence of radial flows, analogously to equation (3) we can write the rate of change of metals in the reservoir as

$$\dot{m}_{\text{Z,gas}} = \dot{m}_{\text{Z,in}} - \dot{m}_{\text{Z,out}} - \dot{m}_{\text{Z,*}} + \dot{m}_{\text{Z,return}}. \quad (6)$$

This contains two source terms and two sink terms. The components are as follows.

- (i) $\dot{m}_{Z,\text{in}}$ represents the metals introduced from the metal-poor halo. We define this gas to have a typical metallicity Z_0 .
- (ii) $\dot{m}_{Z,\text{out}}$ is the metal mass entrained in wind driven outflows. The metallicity of this gas is that of the reservoir, Z .
- (iii) $\dot{m}_{Z,*}$ represents the mass locked into long-lived stars, removing gas with metallicity Z .
- (iv) $\dot{m}_{Z,\text{return}}$ is the metal enrichment resulting from star formation. The characteristic yield, y , is defined as the metal mass returned per unit mass in long-lived stars.

With these principle assumptions, equation (6) can be expressed as

$$\dot{m}_{Z,\text{gas}} = Z_0\dot{m}_{\text{in}} - Z(1 - R + \lambda)\text{SFR} + y(1 - R)\text{SFR}. \quad (7)$$

The rate of change of reservoir metallicity can be written

$$\dot{Z} = \frac{1}{m_{\text{gas}}}(\dot{m}_{Z,\text{gas}} - Z\dot{m}_{\text{gas}}). \quad (8)$$

L13 find that the metallicity of such a system will approach equilibrium on a time-scale shorter than the depletion time-scale (i.e. $\leq t_{\text{dep}} = \epsilon^{-1}$). In which case they show the equilibrium metallicity to be

$$Z_{\text{eq}} = Z_0 + \frac{y}{1 + r_{\text{gas}} + (1 - R)^{-1} \left(\lambda + \epsilon^{-1} \frac{d \ln(r_{\text{gas}})}{dr} \right)}. \quad (9)$$

We now have an expression for the equilibrium metallicity of the system as a function of r_{gas} , a quantity we have already obtained (see Section 5.1). We highlight that there are other conceptually interesting ways of interpreting r_{gas} , which can be alternatively be written as

$$r_{\text{gas}} = \frac{m_{\text{gas}}}{m_*} = \frac{\epsilon^{-1}\text{SFR}}{m_*} = \epsilon^{-1}\text{sSFR}, \quad (10)$$

where sSFR is the specific star formation rate.

By fitting their model to star-forming galaxies from the SDSS, L13 estimate $\epsilon^{-1} \frac{d \ln(r_{\text{gas}})}{dr} \approx -0.25$, and we shall adopt this value. We note, however, that the model is not strongly sensitive to this factor, owing to the degeneracies arising from fitting the y and Z_0 parameters (Section 6.1.4).

6.1.3 Estimating the mass-loading factor, λ

We have parametrized wind-driven outflows via $\dot{m}_{\text{out}} = \lambda \cdot \text{SFR}$, where λ is a mass-loading factor. We shall consider two scenarios, one with winds and the other without. We note here that this choice is not important, for we will show in Section 6.2 that these two scenarios are highly degenerate with y and Z_0 parameters, which we shall fit.

In the simple windless scenario, we will set $\lambda = 0$ everywhere.

For our more complex windy model, we consider outflows that are powered by momentum-driven winds from supernovae (SNe). We follow the prescription described in Dutton & van den Bosch (2009) to define the mass-loading factor

$$\lambda = \frac{p_{\text{SN}}\eta_{\text{SN}}}{V_{\text{esc}}(\mathbf{r})}, \quad (11)$$

where $p_{\text{SN}} = 3 \times 10^4 M_{\odot} \text{ km s}^{-1}$ is the momentum per SN, $\eta_{\text{SN}} = 8.3 \times 10^{-3}$ is the number of SNe per solar mass of stars formed, and $V_{\text{esc}}(\mathbf{r})$ is the escape velocity at a given point, \mathbf{r} , in the disc. The escape velocity itself is defined in terms of the gravitational potential

$$V_{\text{esc}}(\mathbf{r}) = \sqrt{2|\Phi_{\text{tot}}(\mathbf{r})|}, \quad (12)$$

where the gravitational potential, Φ_{tot} , is the sum of contributions from stars, gas and dark matter

$$\Phi_{\text{tot}}(\mathbf{r}) = \Phi_*(\mathbf{r}) + \Phi_{\text{gas}}(\mathbf{r}) + \Phi_{\text{DM}}(\mathbf{r}). \quad (13)$$

We must calculate the gravitational potential at the position in the galaxy where each reservoir/zone is situated. To achieve this, we model the three potentials separately. We shall use thin discs to represent both the stellar and gaseous components. The dark matter component we assume to be distributed in a spherically symmetric halo.

We determine the stellar contribution using the stellar mass maps (Section 2.3). Assuming that the stars lie in a thin plane, we assign a point mass to every map pixel. The potential at any point is the galaxy is then calculated as a sum of the individual point mass potentials, i.e.

$$\Phi_*(\mathbf{r}) = - \sum_i \frac{G \Delta m_i}{|\mathbf{r}_i - \mathbf{r}|}, \quad (14)$$

where G is the gravitational constant, Δm_i is the mass of a pixel, and $|\mathbf{r}_i - \mathbf{r}|$ is the distance in the plane of the galaxy to the centre of the mass pixel.

To estimate the gravitational potential arising from the gas, we adopt the following characteristic surface density profile from Bigiel & Blitz (2012)

$$\frac{\Sigma_{\text{gas}}}{14 M_{\odot} \text{ pc}^{-2}} = 2.1 \exp(-1.65r/r_{25}), \quad (15)$$

where r_{25} is the optical radius where the surface brightness becomes 25 mag arcsec⁻². We use equation 2.164a from Binney & Tremaine (2008) to calculate the gas contribution to the potential. It was shown in Wang et al. (2014) that the characteristic surface density profile provides a good description of the H I-rich galaxies in our sample. Admittedly the profile does not provide as good an approximation to the control sample. However, since the contribution of the gas to the total potential is small ($\lesssim 10$ per cent) this will not affect our conclusions.

The dark matter halo provides the dominant contribution to the halo, making up 50–80 per cent of the total potential, but it is also the most uncertain as we do not have direct constraints on its properties. In view of this, we follow common practice to parametrize the dark matter halo mass distribution with the spherically symmetric NFW profile (Navarro, Frenk & White 1997). In order to do so we need to estimate the halo mass and concentration. We get the former from halo mass–stellar relation derived by Guo et al. (2010, their equation 3) and the halo concentration from Macciò et al. (2007, their fig. 3) with the virial radius of the halo using equation 3 from Dutton & van den Bosch (2009). With this best guess dark matter potential, combined with the potentials of the stellar and gas discs, we are now able to estimate the mass-loading factor λ .

6.1.4 Fitting y and Z_0

Two components in the model remain unconstrained, namely the yield and the metallicity of the infalling gas.

The stellar yield, y , represents the metallicity of the gas returned by short-lived stars. If we assume there is a universal IMF, then we expect y to be constant between galaxies and independent of location within a specific galaxy. The stellar yield can in principle be calculated from stellar evolution models. However, the large and poorly understood systematic offsets between the various gas-phase metallicity indicators (see Appendix A) mean that we are unable to

determine absolute abundances for our galaxies so we have decided to assume that y is the same for all galaxies, but unknown so we fit it as a global constant.

The metallicity of gas infalling from the halo, Z_0 , is a poorly known quantity. For simplicity, we therefore assume that the infalling gas has the same metallicity at all radii for each galaxy and that the haloes of all the Bluedisk galaxies have the same metallicity. This may not be a bad assumption since all the galaxies are of similar total stellar mass, and therefore may possess similar mass haloes. As with the stellar yield, our prediction of Z_0 also suffers from effects of systematic offsets due to the choice in metallicity indicators. Thus, we also make Z_0 a global constant that is to be fit.

In summary, given the stellar mass maps and gas mass distributions, the resulting model, equation (9), has two free parameters, y and Z_0 . These global parameters, namely the stellar yield and the halo metallicity, are fit for all radial bins, across all galaxies, simultaneously. Due to the systematic offsets between metallicity indicators, we caution that inferences should not be made on the fitted values themselves.

6.2 Bluedisk galaxies as local gas regulators

Having outlined the local gas regulator model, we demonstrate the results for all 50 Bluedisk galaxies in Fig. 12. Surprisingly this simplistic equilibrium model appears to match well for many, but by no means all, of the galaxies. It is strikingly clear that with only two globally-set free parameters, we can reproduce a large variety of observed metallicity profile shapes that these galaxies exhibit. The model also reproduces the observed outer metallicity drops, which is attributed to the transition from a stellar-dominated inner disc to a more gas-dominated outer disc.

The centres of the galaxies appear to be most problematic for the local gas regulator model to reproduce. One of our key model assumptions is that we assume independence between radial zones. But the presence of bars and bulges at the centres of galaxies might invalidate this assumption. For example, a bar could be expected to drive strong radial flows inwards, which if this were the case, we could expect steepened metallicity gradients (Goetz & Koeppen 1992). However, we see no obvious connection between deviations from the local gas regulator model and the presence of a strong bars or a prominent bulges.

Alternatively, the failure of the model may indicate that some of these deviant central regions are contaminated by emission whose origin is not photoionization, such as shocks and LINERs, which our selection criteria have failed to exclude. Emission from non-photoionizing origins can impact different metallicity determination methods in different ways. Although as shown in Appendix A different metallicity indicators yield different absolute and relative abundances, we should expect there to exist a monotonic mapping between the methods. For example, in Fig. 7 we can see the inner regions of BDs 2, 15 that they show contradictory behaviour of the metallicity of the CL01 and KK04 indicators. This primarily indicates contamination from non-photoionizing origins, thus it is not surprising the local gas regulator model appears to fail in these regions.

We find no difference in the quality of fit of the local gas regulator model between the H I-rich and control galaxies. Overall both our H I-rich, control galaxies fitting equally well. However, our local gas regulator model should not be expected to succeed for galaxies that are interacting, since interactions could also drive strong gas flows. Indeed, our excluded (non-isolated) sample of galaxies exhibit some

of the most discrepant fits, e.g. BD 31, which has a very different metallicity profile from the one predicted by the model.

We construct the local gas regulator model with momentum-driven winds and windless cases, both of which appear to represent the data more or less equally well. This is not because winds have no effect, indeed they do impact strongly on the metallicity, but the effects of the wind is largely degenerate with our fitted parameters: y , and Z_0 . The wind model we employ modifies the metallicity changing the peak metallicity and inner slope (where $r_{\text{gas}} \sim 0$). The loss of metals due to winds can be compensated by increasing the values of y and/or Z_0 . Since we are forbidden from interpreting the values of y and Z_0 , we are unable to conclude anything either for or against the existence of enriched wind-driven outflows.

We note that our results here are compatible with the recent work of Ho et al. (2015). Using an analytical formalism similar to the L13 gas regulator they are able to reproduce the distribution of metallicity gradients observed. Therein the metallicity profile is also determined by the current ratio of gas-to-stellar mass.

7 CONCLUSIONS

We present radial gas-phase metallicity profiles of 50 late-type galaxies that form the Bluedisk survey. We explore how the H I content of these galaxies affects their metallicity profiles. Although we find a correlation between H I mass fraction and the metallicity gradient, we observe that the metallicity profiles of our H I-rich and control samples show remarkable similarity. Furthermore we find that using a simple equilibrium model we are able to approximate the metallicity profiles of both samples with equal success. We summarize our main results as follows.

- (i) We confirm the local mass–metallicity relation for the Bluedisk galaxies. Although we note that at low stellar mass density there appears to be a residual anticorrelation of metallicity with gas mass density.
- (ii) The metallicity gradient of the galaxies is strongly correlated with their H I mass fraction. Galaxies with higher H I mass fractions have steeper metallicity gradients.
- (iii) We find that in some galaxies the outer disc exhibits steeper metallicity gradients than in the inner disc. However, unlike previous work that has shown this, we find these galaxies to be ubiquitous in both the H I-rich and control samples.
- (iv) The barred galaxies in our sample bars tend to have flatter metallicity profiles. This is not sufficient, however, to explain observed differences between the metallicity profiles of the H I-rich and control samples.
- (v) By applying a simple equilibrium analytical model, we find that we are able to approximate the metallicity profile shapes with the ratio of gas-to-stellar mass, $r_{\text{gas}} = \Sigma_{\text{gas}}/\Sigma_{*}$. In the outer disc where Σ_{*} is low, a transition to $r_{\text{gas}} > 1$ occurs. This naturally gives rise to the steeper outer metallicity gradients which are observed.

If metallicity is truly in equilibrium, this would naturally explain the local mass–metallicity relation. Also, since the dynamic range of the stellar mass density is much greater than the gas mass density, the overall metallicity profile represents the integrated build up of metals.

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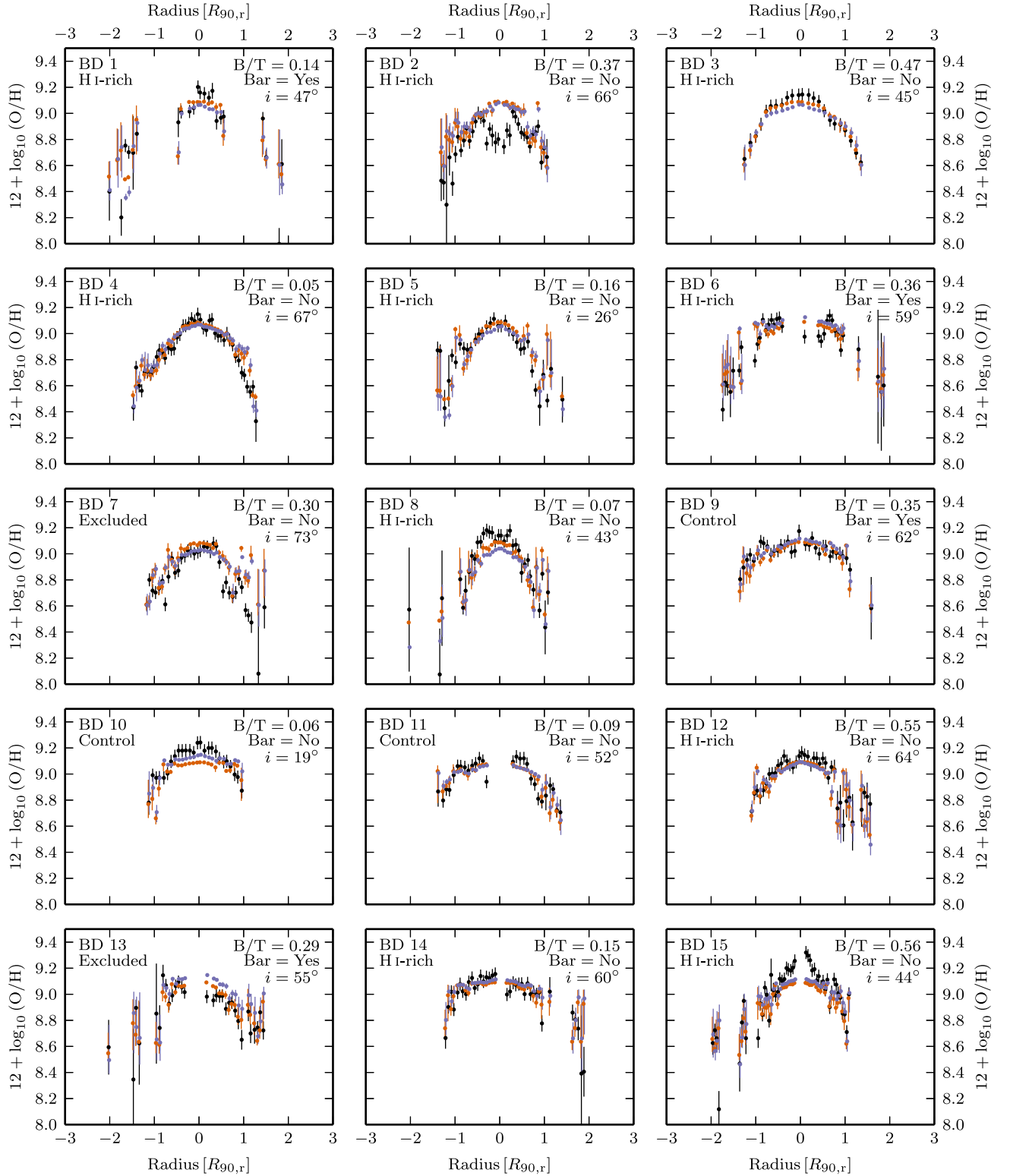
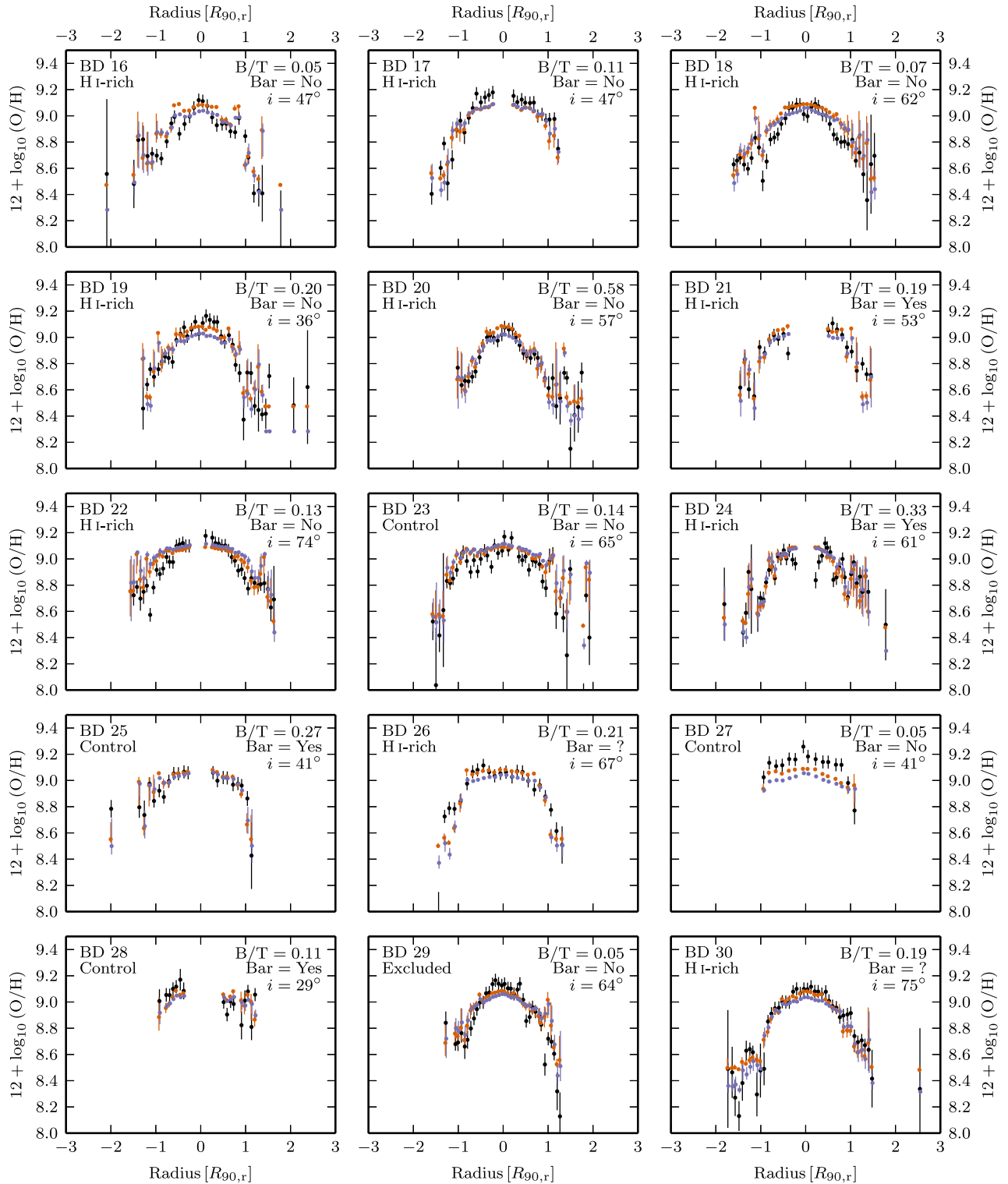


Figure 12. The local gas regulator model compared to the metallicity profiles of the Bluedisk galaxies. The observed CL01 metallicity is shown in black. Windless and windy models are shown in orange and blue, respectively. Error bars on the models are not the true full error, but rather they indicate the effects of $\pm 1\sigma$ deviations in the gas-to-stellar mass ratio, r_{gas} . In the top right of each plot, we label the bulge-to-total light ratio, B/T. We also denote whether a bar is present. In highly inclined systems where we would not be able to determine the presence of a bar, we denote this with a “?” symbol. Since our modelling may be problematic at high inclinations, we also include the measured inclination i . The globally fitted parameter values are $y = 1.27 \times 10^{-3}$, $Z_0 = 1.91 \times 10^{-4}$ and $y = 0.55 \times 10^{-3}$, $Z_0 = 2.97 \times 10^{-4}$ in the windy and windless cases, respectively.

Figure 12 – *continued.*

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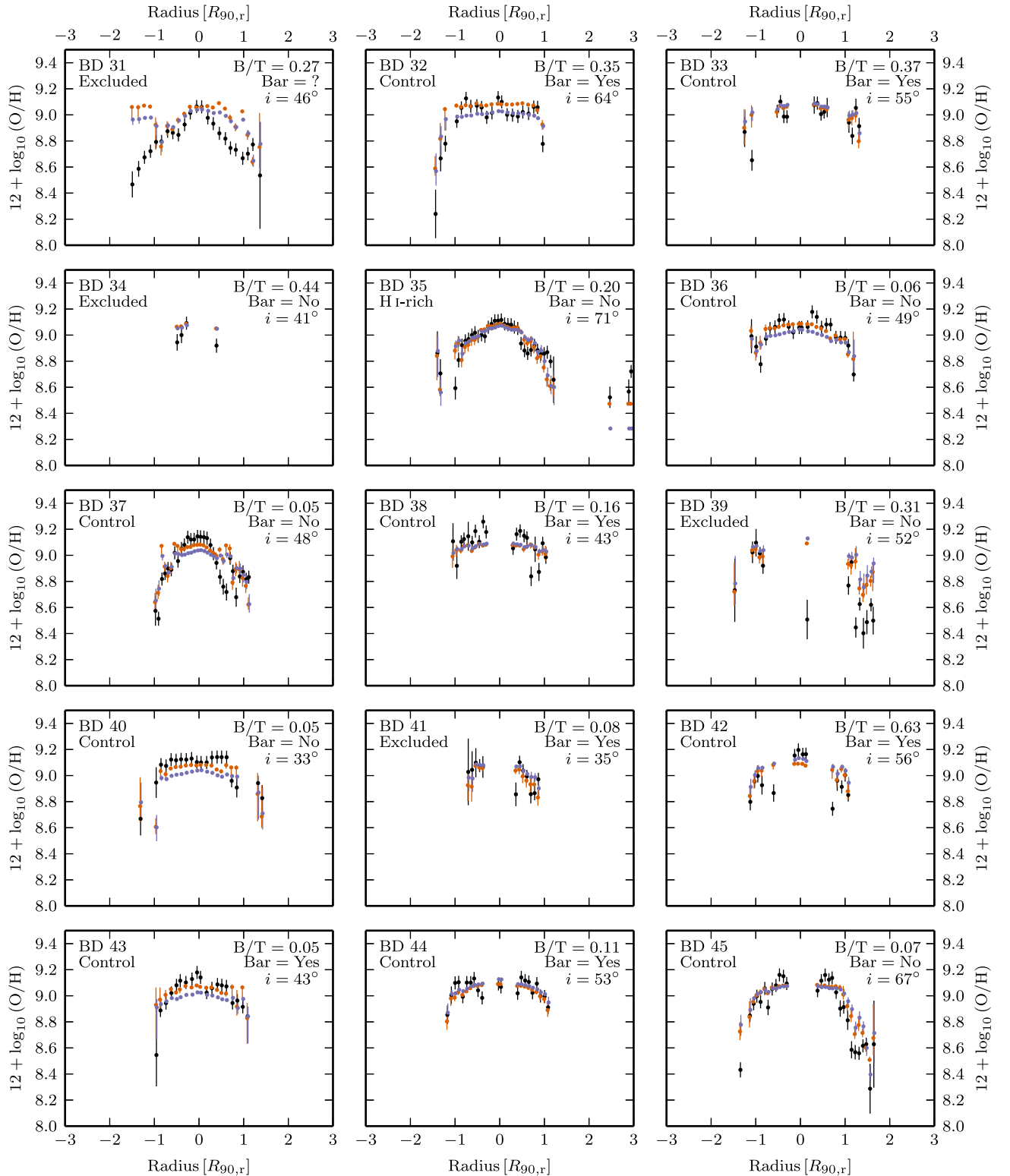
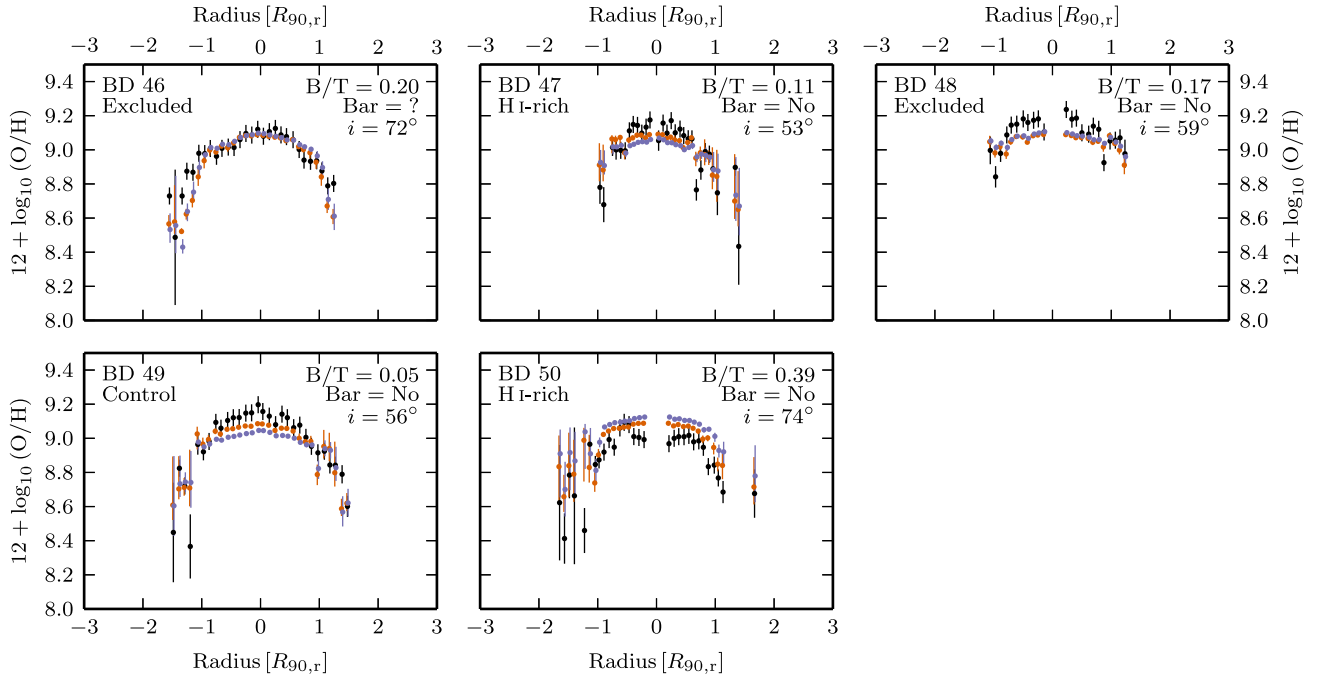


Figure 12 – continued.

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Figure 12 – *continued.*

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APPENDIX A: COMPARISON OF METALLICITY INDICATORS

In Section 4.3, we discuss a variety of methods for determining metallicity. We shall now discuss the similarities and differences between these methods and justify the use of our primary method that uses the CL01 models.

In Figs 7, 8 and 9, we present metallicities derived from the CL01, KK04 and NS methods. It should be immediately apparent that large (0.6 dex) systematic offsets in metallicity exist between these methods. The CL01 and KK04 methods, both derived from theory, produce largely consistent metallicities. However, they report much higher metallicities than those from the NS method, which is derived empirically. It is also noticeable that the NS method produces shallower profiles than the CL01 models. Of course disparities between metallicity indicators are not unique to this work and they have been well documented by Kewley & Ellison (2008) and Moustakas et al. (2010).

Although it is not desirable to have absolute and relative differences between metallicity methods, we assert that this is of no major significance provided the metallicity indicators are

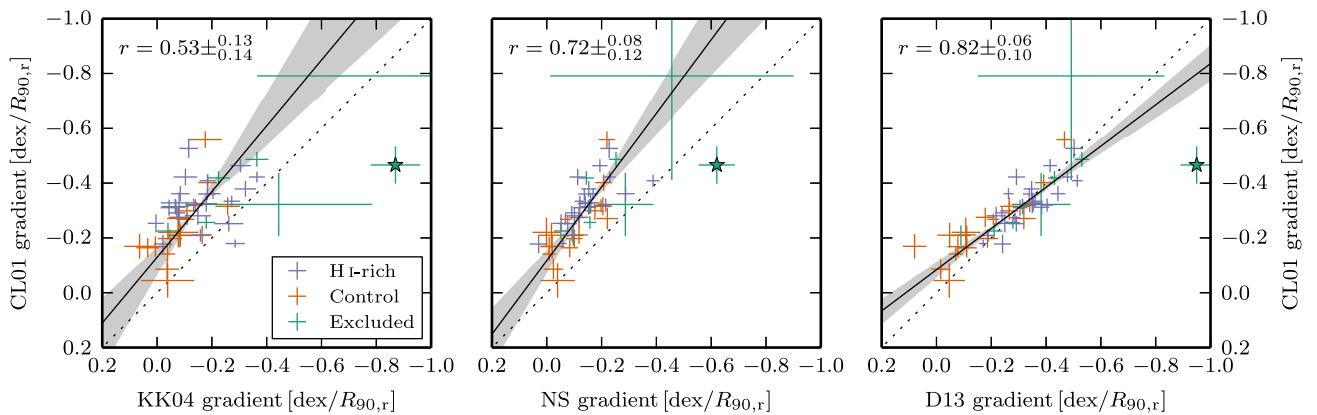


Figure A1. Comparison of straight-line gradients derived for CL01 against three difference metallicity indicators (KK04, NS and D13), for all 50 Bluedisk galaxies, H I-rich (blue), control (orange) and excluded (green). Linear fit, using orthogonal-distance regression, is shown (solid line) and the shaded region indicates the associated error in the slope. The dotted line indicates equal $x = y$ mapping. The Spearman's rank correlation coefficient is given in the top-left corner of each figure. A star indicates BD 39, which is excluded from the regression and r -statistic computation, due to its companion galaxy. Median values and $\pm 1\sigma$ errors on statistics and gradients are computed by bootstrapping Monte Carlo-scattered resampled data. For reference, $r = 0.34$ is the one-tailed Spearman's r -value at a $\alpha = 0.01$ significance level.

consistent. In other words, a galaxy with a steeper profile in one method should yield a steeper profile in all methods. We test this assumption by fitting simple straight-line models to metallicity profiles derived from different methods, and compare the inferred gradients. In Fig. A1, we show the gradients inferred from the CL01 method versus KK04, NS and D13 methods. We note the strongest correlation exists between the CL01 and D13 models, indicating our results are not strongly dependant on the photoionization models employed. However, particularly remarkable is the similarity between CL01 and NS gradients, since the NS method is empirically calibrated. We note that the KK04 method often produces shallower gradients than CL01. Indeed, it appears that the KK04 methods can produce much higher metallicities than CL01 in the outskirts of the galaxies (e.g. BD 50). This is often associated with significant differences in inferred dust attenuation strengths, with the standard Case B method, producing erratic results between adjacent bins. Nevertheless, we derive comfort in using our CL01 metallicities as our default method for metallicity determination.

APPENDIX B: GAS SURFACE DENSITY ESTIMATES

In Section 4.4, we use a technique developed by B13 to estimate gas surface mass densities from optical lines. It was shown there that when most of the lines in the optical spectrum are available it is possible to use photoionization models with a flexible treatment of metal depletion to place constraints on the gas surface mass density of galaxies.

This works because emission line ratios are sensitive to temperature and since metals are very important coolants, changing their depletion factor at fixed metallicity, Z , changes the temperature in the gas noticeably. Exploiting this fact, B13 showed that it is possible to place constraints on the dust-to-metal ratio, ξ , of ionized gas

using only optical emission lines. When combined with an estimate of metallicity and the dust optical depth, primarily from Balmer lines, and a simple model for the ISM of a galaxy they show that total gas surface mass densities can be estimated through

$$\Sigma_{\text{gas}} = 0.2 \frac{\tau_V}{\xi Z} M_{\odot} \text{pc}^{-2}. \quad (\text{B1})$$

We have calculated this quantity for each spatial bin in the spectra discussed in this paper.

We shall now provide an additional check of these spectroscopic Σ_{gas} estimates. In Wang et al. (2014), herein referred to as Paper II, azimuthally averaged Σ_{gas} were calculated. These were computed by combining the observed H I surface density with an estimated contribution from H₂. The H₂ component was estimated using an SF scaling relation applied to the observed SFRs. In Fig. B1, we show a few select examples of our spectroscopic gas profile against those from Paper II. These galaxies were selected to span a range from very poor to very good agreement. Overall the match between the two estimators is reasonable given the differences in analysis and that the profiles from Paper II are azimuthally averaged while the spectrally determined gas surface densities originate from long-slit spectra.

Nevertheless, it is noticeable that the spectroscopically determined gas surface densities do not show a strong drop in the outer regions of the galaxies. This is possibly due to a characteristic of the B13 method that is not discussed in detail by B13 (but see their fig. 15), namely that it might give a biased estimate of the *average* gas surface density in the outer regions of galaxies. The reason for this is that the method only works reliably when there is a clear emission line source, in practice an H II region in the spectral aperture. In the outskirts of galaxies these regions are fewer and tend to coincide with peaks in the local gas density. But these peaks provide biased estimates of the azimuthally averaged gas surface density at those radii so the spectroscopic method will also provide biased estimates.

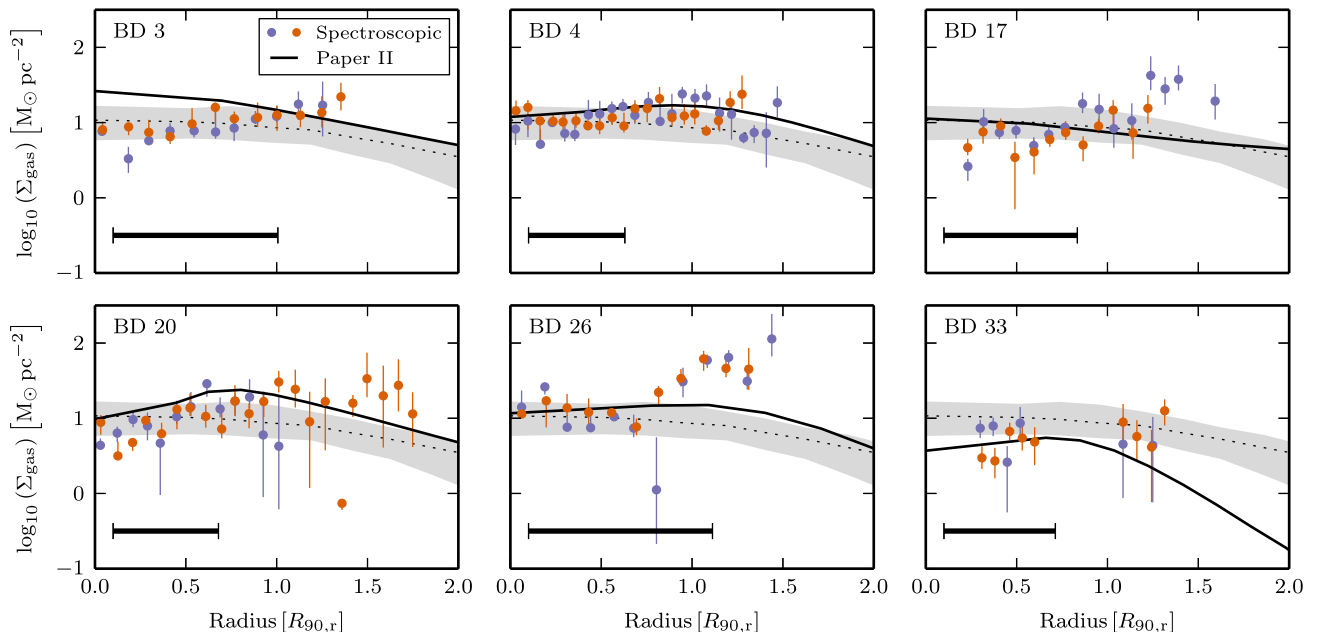


Figure B1. Comparison of gas density profiles of six galaxies using two different estimators. The spectroscopic Σ_{gas} estimates are represented by blue and orange data points, with the two colours distinguishing either side of the spectroscopic slit. The azimuthally averaged Σ_{gas} estimates from Paper II are shown as a solid black line. We also show the Σ_{gas} profile averaged over all galaxies presented in Paper II. These are plotted as a dotted black line and a grey shaded area indicating the median and $\pm 1\sigma$ quantile range, respectively. In the bottom-left corner of each plot, a thick black bar indicates the scale of 13 arcsec which roughly approximates the resolution limit of the Paper II estimates.

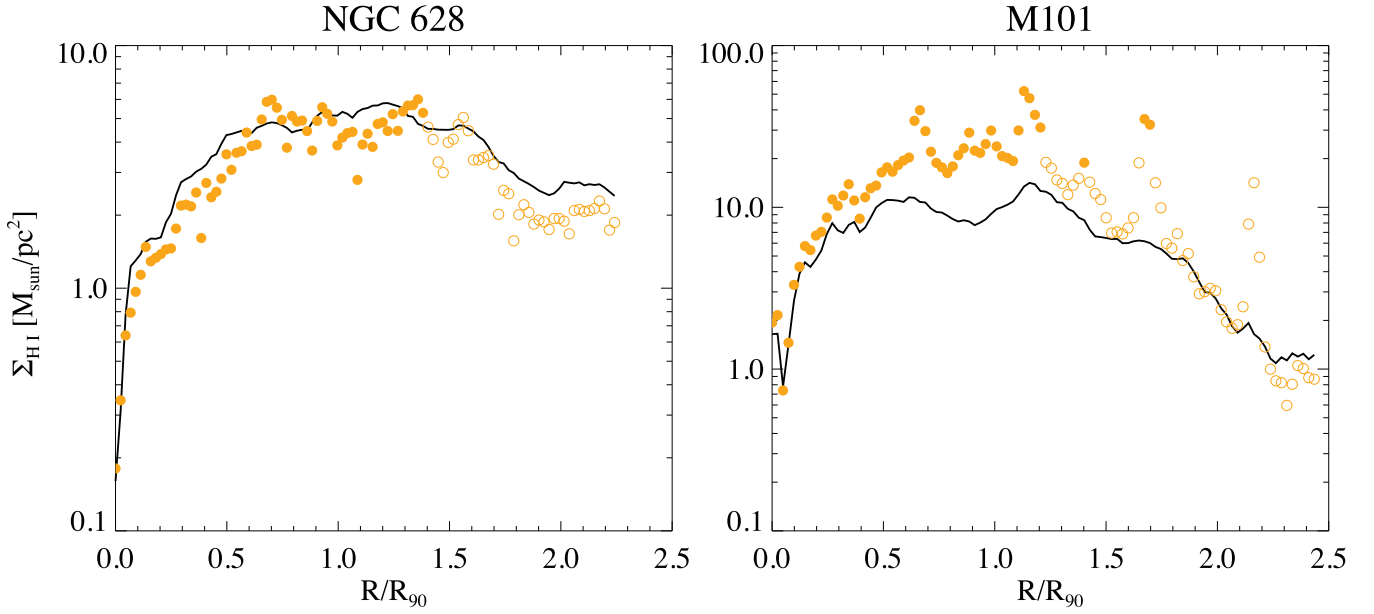


Figure B2. The azimuthally averaged H I gas density in galaxy discs (solid black line), compared to the H α flux weighted density (orange symbols). The left-hand panel shows the results for NGC 628 with the right-hand panel that of M 101. The open orange symbols indicate annuli where the average SFR is $< 10^{-3} \text{ M}_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}$.

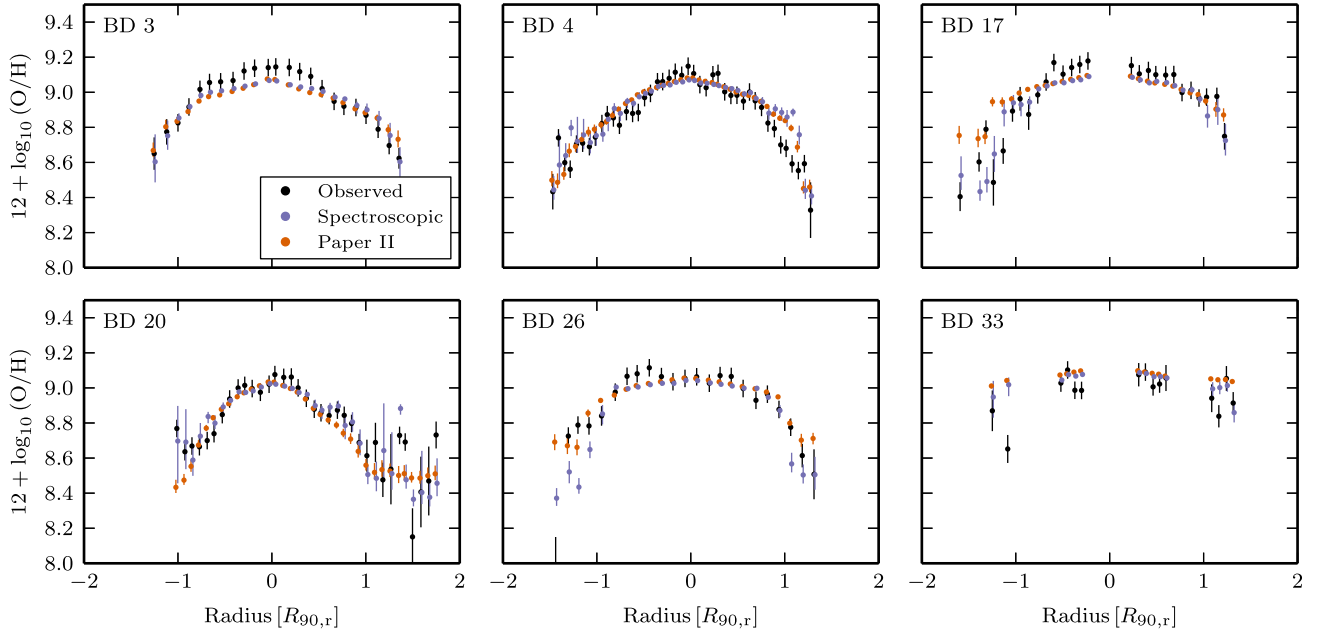


Figure B3. Comparing the effects of two gas density estimators on the local gas regulator model. The observed CL01 metallicity is shown in black. The model using spectroscopic Σ_{gas} estimates is shown in blue. The equivalent model using Paper II Σ_{gas} estimates is shown in orange. Both models assume the modelling scenario with winds.

To illustrate this fact, Fig. B2 shows the azimuthally averaged H I gas profiles for two large nearby spiral galaxies, NGC 628 and M 101 as black solid lines. The H I maps were taken from Walter et al. (2008). To illustrate the effect of probing the gas density at the location of H II regions in the outer discs, we overplot the H α weighted H I profiles as orange symbols on top. The open symbols are for annuli where the mean SFR is $< 10^{-3} \text{ M}_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}$, assuming a Salpeter IMF and the $L(\text{H}\alpha)$ to SFR conversion factor of (Kennicutt 1998). We took the H α maps from Hoopes, Walterbos & Bothun (2001).

What is noticeable is that the two galaxies are rather different with the H α weighted profile in NGC 628 being very close to the straight mean profile. In contrast, the H α -weighted profile in M101 is noticeably higher than the mean profile and it is also clear that star formation at large radii is connected to relatively high gas densities.

Nevertheless, the differences between two estimators of Σ_{gas} do not strongly impact on our analysis using the local gas regulator model (Section 6.2). This is demonstrated in Fig. B3, which highlights the relatively minor effect of the choice of gas-density estimator on our conclusions. The reason for the lack of significant

difference is that the local gas regulator models depends on the ratio of gas-to-stellar mass. Across a galaxy the dynamic range of the stellar mass density is much greater than that of the gas-mass density. Therefore the overall shape of the metallicity profile is primarily set by the stellar mass density profile.

APPENDIX C: METALLICITY PROFILES EQUALLY WEIGHED BY GALAXY

In Section 5.2.1, we produce stacked metallicity profiles. As each data point receives an equal weight within a bin. Galaxies with many data points may, however, dominate a bin. If these galaxies have

atypical metallicity profiles, this would be mimicked in the stacked profile. Such problems are only likely to arise in the outermost radial bin of a stack. To qualify and counteract this effect we also produce stacked profiles we use a weighted median, where data points are each weighted inversely to the number of data points from the same galaxy per bin.

In Fig. C1, we reproduce Fig. 11 with this new weighting. The only appreciable difference occurs for some control sample stacks, where the outermost metallicity is never lower than the metallicity of the H I-rich galaxies. Since these bins are dominated by a few galaxies, they may not be representative of the whole sample.

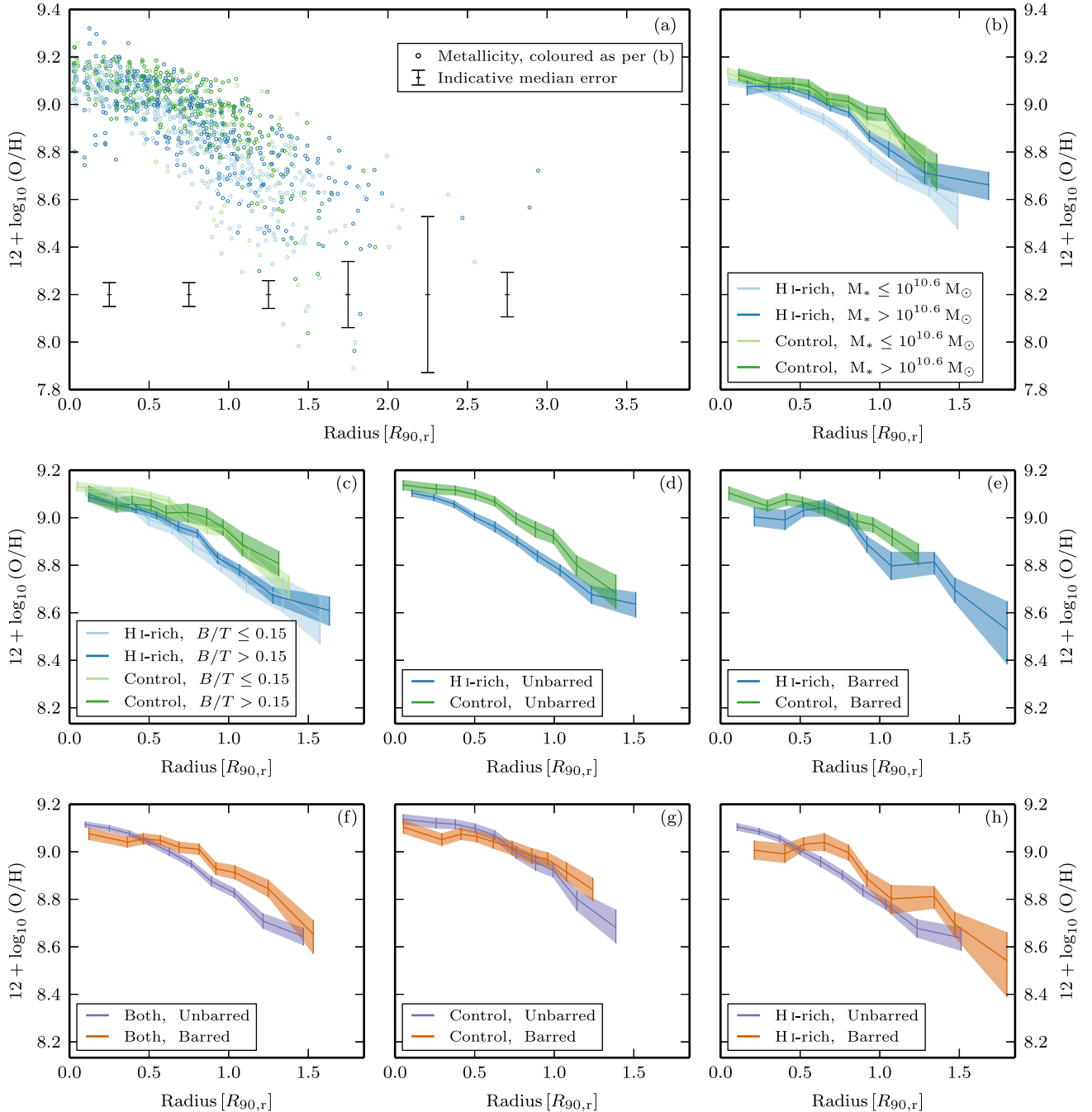


Figure C1. Equivalent of Fig. 11 with weighted such that galaxies receive equal weight per bin. See Fig. 11 for description.

Table D1. Positions and orientations of the spectroscopic slit observations: Galaxy identification number, Common name, Right ascension, Declination, and Position angle of slit. PA is defined with (North = 0°, East = 90°)

ID	Name	RA (deg)	Dec. (deg)	PA (deg)
BD 1	UGC 4283	123.591750	+39.251361	133
BD 2	UGC 4429	127.194792	+40.665889	32
BD 3	MCG+07-18-029	129.277167	+41.456333	3
BD 4	IC 2387	129.641667	+30.798694	20
BD 5	NVSS J084916+360710	132.318292	+36.119806	70
BD 6	NGC 2668	132.344042	+36.710333	160
BD 7	UGC 4615	132.356167	+41.771250	15
BD 8	UGC 4798	137.177583	+44.810667	140
BD 9	NVSS J091458+512138	138.743000	+51.361056	18
BD 10	NGC 2895	143.104375	+57.482889	40
BD 11	NGC 3135	152.726583	+45.950361	84
BD 12	UGC 5534	154.042708	+58.427000	98
BD 13	UGC 6183	166.995917	+35.463278	5
BD 14	UGC 6755	176.739750	+50.702139	66
BD 15	NGC 3897	177.247750	+35.016056	135
BD 16	NVSS J125203+514050	193.014792	+51.680056	103
BD 17	UGC 8205	196.806625	+58.135028	107
BD 18	UGC 8338	199.015042	+35.043528	82
BD 19	NGC 5497	212.631833	+38.893556	49
BD 20	UGC 9429	219.499750	+40.106194	112
BD 21	2MASX J16073420+3629026	241.892583	+36.484028	109
BD 22	UGC 10523	250.793500	+42.192778	173
BD 23	UGC 10553	251.811625	+40.245083	145
BD 24	MCG+10-24-123	259.156042	+58.411889	169
BD 25	MCG+10-25-046	262.156333	+57.145056	92
BD 26	2MASX J07274518+4210499	111.938042	+42.180722	50
BD 27	2MASX J08024061+3431171	120.669375	+34.521444	165
BD 28	MCG+09-14-017	123.309125	+52.458722	150
BD 29	UGC 4427	127.312167	+55.523000	180
BD 30	UGC 4863	138.603167	+40.777917	169
BD 31	MCG+08-17-066	139.191875	+45.813735	49
BD 32	2MASX J0918351+321611	139.646250	+32.270000	57
BD 33	UGC 5016	141.539292	+49.310194	32
BD 34	NGC 3013	147.539000	+33.569333	70
BD 35	UGC 5346	149.420208	+45.258667	29
BD 36	2MASX J09574902+5149162	149.454542	+51.821194	107
BD 37	NGC 3164	153.797625	+56.672083	3
BD 38	2MASX J10154226+5540030	153.926083	+55.667500	26
BD 39	UGC 5936	162.530375	+36.341833	134
BD 40	MCG+06-25-025	168.563542	+34.154389	18
BD 41	MCG+08-24-070	197.879167	+46.341778	68
BD 42	MCG+08-24-089	198.236250	+47.456667	34
BD 43	MCG+07-28-032	203.374583	+40.529667	55
BD 44	2MASX J13410027+4225525	205.251042	+42.431417	133
BD 45	IC 1074	222.988833	+51.264889	115
BD 46	2MFGC 12932	241.528958	+35.981444	144
BD 47	UGC 10312	244.382542	+31.194472	112
BD 48	NGC 6145	246.259833	+40.946639	4
BD 49	2MASX J17143602+3044011	258.650208	+30.733528	160
BD 50	UGC 10863	261.557250	+62.149472	102

APPENDIX D: THE ATLAS

In Table D1, we provide the positions and orientations of each spectroscopic slit.

In the online edition of this paper, we provide an atlas of the Bluedisk galaxies and observed data in Figs D1–D50. Each contains four panels, the x -axes of which are aligned.

The first panel shows an SDSS gri composite image of the galaxy. Two horizontal dashed lines indicate orientation of the 3 arcsec

spectroscopic slit. The second panel shows the metallicity profiles as shown in Figs 7–9. In the third panel, we indicate the flux of four emission lines ($H\alpha$, $H\beta$, $N\text{II}$ and $O\text{III}$). In the fourth panel, we plot the inferred velocity shifts of the emission lines from the red-arm of the spectrograph. These velocities have been corrected for the effect of inclination.

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